

# Structural, Stabilities, and Electronic Properties of Bimetallic Mg<sub>2</sub>-doped Silicon Clusters

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The equilibrium geometries, relative stabilities, growth patterns, and electronic properties of magnesium-doped silicon clusters Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1–11$ ) have been systematically investigated at the B3LYP/6-311G (d) level. A large number of initial configurations are optimized and the lowest-energy stable geometries of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1–11$ ) clusters with different spin multiplicities are determined. The results indicate that the most stable configurations for Mg<sub>2</sub>Si<sub>n</sub> clusters favor the three-dimensional structures at  $n = 3–11$ . The analyses of the averaged binding energies, fragmentation energies, second-order energy difference, and HOMO–LUMO gaps suggest that the Mg<sub>2</sub>Si<sub>4</sub> and Mg<sub>2</sub>Si<sub>6</sub> clusters have the stronger relative stability, and magnesium atoms doping enhances the chemical activity of the silicon framework. The natural population and natural electronic configuration analyses show that the charge transfer occurs from the 3s orbital of the magnesium atoms to the silicon atoms and 3p orbital of the magnesium atoms.

*Key words:* Mg<sub>2</sub>Si<sub>n</sub> Clusters; Density Functional Theory; Geometrical Structures; Electronic Properties.

## 1. Introduction

Alkaline-earth metal silicides have been the subject of considerable interest for several decades, not only because they provide a model system for the detailed study of silicon-based microelectronics, but because of the undeniable importance of silicon in the semiconductor industry [1–3]. They have unusual physical and chemical properties, such as low resistivity, high elevated temperature strength, high specific strength, high elastic modulus, and low thermal conductivity, and wide applications in many fields, which include the usage in the aviation and automobile industries. However, there is little known about the molecular level mechanism, partly due to the difficulty in studying, both theoretically and experimentally, a typically solid–solid reaction on the molecular scale. Small clusters offer a new way in studying such systems, which can provide insight into the strength and nature of the metal–silicon bonds [4–9]. Beck [10] reported the observation of a reaction between a metal atom and silicon in a supersonic jet to

form metal atom–silicon clusters. Based on the anion photoelectron spectroscopy, Grubisic et al. [11] investigated several LnSi<sub>n</sub> cluster anion systems (Ln=Yb, Eu, Sm, Gd, Ho, Pr), which expands the range of studied lanthanide–silicon systems to encompass half of the lanthanide series. Kishi et al. [12] performed the photoelectron spectroscopy measurements on the silicon–sodium cluster anions. The ionization energy and electron affinities of Si<sub>n</sub>Na<sub>m</sub> revealed several patterns with the number of the sodium atoms, and the charge transfer mechanism in the cluster reflects the nature of the electron-accepting orbitals of the corresponding Si<sub>n</sub> and the adsorption types of the sodium atoms. To explore the structural evolution of chromium-doped silicon clusters, Kong et al. [13] investigated CrSi<sub>n</sub><sup>−</sup> ( $n = 3–12$ ) clusters using the anion photoelectron spectroscopy and density functional theory (DFT) calculations. The combination of experimental measurement and theoretical calculations suggested that the onset of endohedral structure in CrSi<sub>n</sub><sup>−</sup> clusters occurs at  $n = 10$ , and the magnetic properties of the CrSi<sub>n</sub><sup>−</sup> clusters are correlated with their

geometric structures. By using first-principles DFT-generalized gradient approximation (GGA) calculations, Guo et al. [14] reported the stabilities and magnetic properties of the transition metal encapsulated MSi<sub>n</sub> (M=Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn;  $n = 8 - 16$ ) clusters.

Although silicon clusters doped with many different metals have been studied extensively in recent years, there are no systematic theoretical and experimental investigations on Mg<sub>2</sub>Si<sub>n</sub> clusters. To gain insight into the geometries and electronic properties of the Mg<sub>2</sub>Si<sub>n</sub> clusters and further enrich our knowledge of silicon-based clusters, we carried out the computational studies on a series of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters. We hope that the theoretical study will be useful to understand the geometric structures of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters in the ground state and can provide useful information in exploring the physical mechanism of the growth behaviors of Mg<sub>2</sub>Si<sub>n</sub> clusters.

## 2. Computational Methods

All geometrical structure optimization and frequency analysis of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters are performed by using DFT with 6-311G (d) basis set, as implemented in the Gaussian 09 program package [15]. The B3LYP exchange-correlation potential [16, 17] is used. To gain the lowest-energy structures, a great deal of initial configurations of Mg<sub>2</sub>Si<sub>n</sub> clusters, which include one-dimensional (1D), two-dimensional (2D), and three-dimensional (3D) configurations, are sought using the following two approaches:

- (i) by considering the numbers of previous reported Si<sub>n</sub>, XSi<sub>n</sub>, and X<sub>2</sub>Si<sub>n</sub> structures [18–23];
- (ii) by placing magnesium atoms at various adsorption or substitutional sites on the basis of optimized Si<sub>n</sub>, XSi<sub>n</sub> or X<sub>2</sub>Si<sub>n</sub> geometries, i. e., Mg-

capped, Mg-substituted, and Mg-encapsulated patterns.

Due to the spin polarization, all configurations are optimized at possible spin multiplicities (singlet, triplet, and quintet) [24]. To ensure the reliability of the theoretical method and basis set used here, the bond length, vibrational frequency, vertical ionization potential (VIP), and vertical electron affinity (VEA) [24] of Si<sub>2</sub>, Mg<sub>2</sub>, and MgSi dimers are calculated and illustrated in Table 1. For the Si<sub>2</sub> dimer, the calculated values are  $r(\text{Si}_2) = 2.17 \text{ \AA}$ ,  $\omega(\text{Si}_2) = 540.37 \text{ cm}^{-1}$ ,  $\text{VIP}(\text{Si}_2) = 9.13 \text{ eV}$ , and  $\text{VEA}(\text{Si}_2) = 2.00 \text{ eV}$ , which are in good agreement with the available experimental results [25–27]. For the Mg<sub>2</sub> dimer, the theoretical results of bond length (3.93 Å) and vibrational frequency (44.91 cm<sup>-1</sup>) are also in good agreement with the experimental results [28]. So, the theoretical method and basis set are reasonable.

## 3. Results and Discussion

### 3.1. Geometrical Structures

For Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters, we have explored a lot of low-lying isomers and determined the lowest-energy isomers. Although many initial structures have been taken into account, only the obtained ground state and some metastable isomers are plotted in Figure 1. According to their energies from low to high, these isomers are designated by *na*, *nb*, *nc*, and *nd*, where *n* represents the number of silicon atoms in the Mg<sub>2</sub>Si<sub>n</sub> clusters.

The lowest-energy structure of Mg<sub>2</sub>Si (1a) is a triplet with C<sub>2v</sub> symmetry and the Mg–Si–Mg apex angle is 72.21°, in which the two magnesium atoms are located on the two sides. The linear isomer Si–Mg–Mg 1b (C<sub>∞v</sub>) with the silicon atom at one end is 0.47 eV higher in energy than the 1a isomer. All possible initial structures of Mg<sub>2</sub>Si<sub>2</sub> clusters with different spin

Table 1. Comparison of the calculated properties for Si<sub>2</sub>, Mg<sub>2</sub>, and MgSi dimers with the corresponding available experimental data.

	Si <sub>2</sub>		Mg <sub>2</sub>		MgSi	
	Cal	Expt	Cal	Expt	Cal	Expt
$r$ (Å)	2.17	2.246 <sup>a</sup>	3.93	3.89 <sup>d</sup>	2.58	–
$\omega$ (cm <sup>-1</sup> )	540.37	511.0 <sup>a</sup>	44.91	45 <sup>d</sup>	288.33	–
VIP (eV)	9.13	> 8.49 <sup>b</sup>	6.59	–	6.82	–
VEA (eV)	2.00	2.176 ± 0.002 <sup>c</sup>	0.04	–	0.59	–

<sup>a</sup>[25], <sup>b</sup>[26], <sup>c</sup>[27], <sup>d</sup>[28]

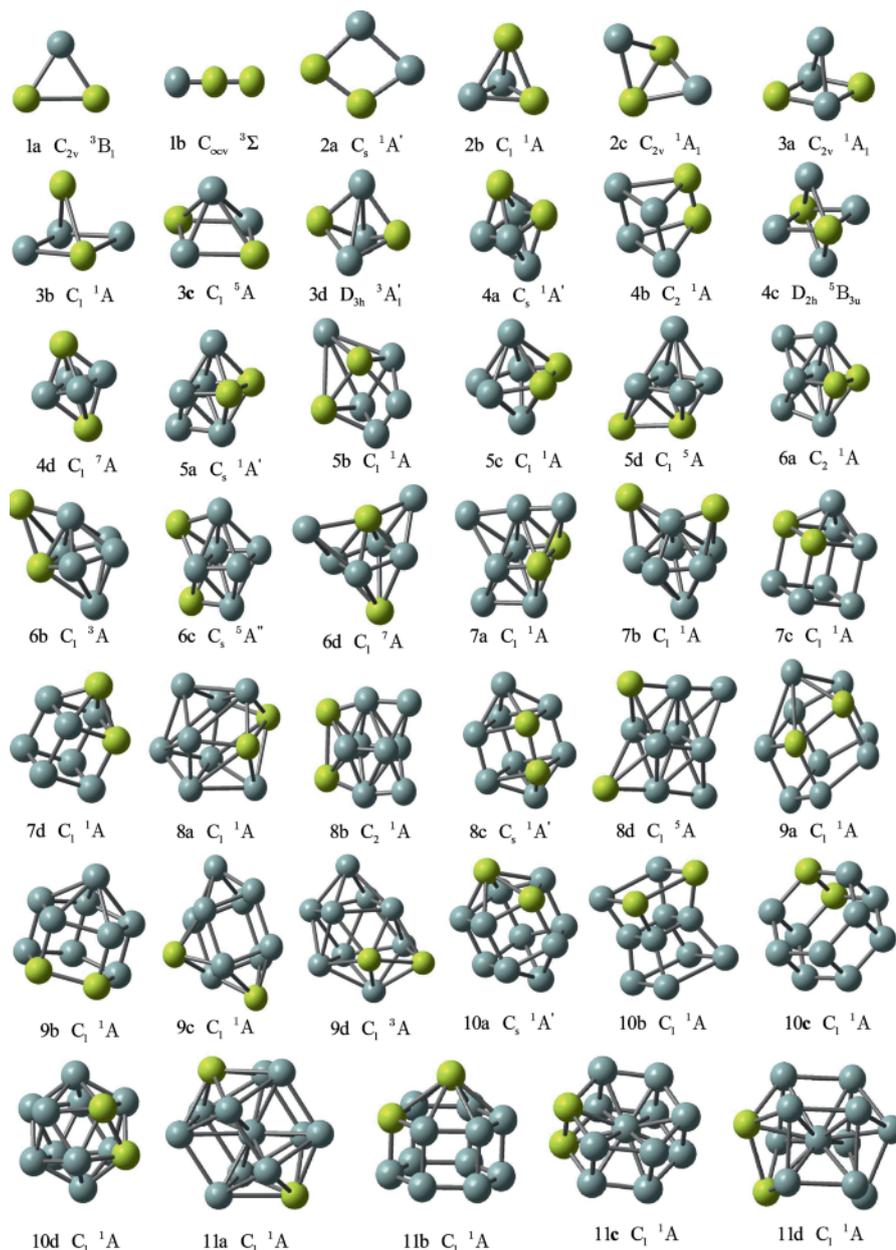


Fig. 1 (colour online). Lowest-energy and low-lying structures of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters.

multiplicities are optimized. The calculated results reveal that the ground state isomer 2a is a rhombus-like structure with singlet state  $^1A'$ , which can be viewed as a planar rhombus lowest-energy structure of Si<sub>4</sub> [18] by replacing two silicon atoms with magnesium atoms. The pyramid isomer 2b with  $^1A$  electronic state is 0.10 eV higher in energy than the 2a isomer. The third

isomer 2c with  $C_{2v}$  symmetry is a butterfly structure with the Mg–Mg bond for the ‘body’ of the insect plus four Mg–Si bonds at the edges of the ‘wings’, and its energy is higher than that of isomer 2b by 0.54 eV. With regard to the Mg<sub>2</sub>Si<sub>3</sub> cluster, the lowest-energy isomer 3a is a pyramid structure with  $C_{2v}$  symmetry, which can be described as a substituted version of the

structure of the ground state Si<sub>5</sub> [18] cluster. The low-lying isomers 3b and 3c are similar to the 3a isomer, the difference between them is the site of magnesium atom, but the energy of the 3a isomer is lower than that of 3b and 3c isomers by 0.65 and 0.77 eV, respectively. From  $n = 3$  afterwards, the 3D structures show higher stability than the planar structures. Among the isomers of Mg<sub>2</sub>Si<sub>4</sub>, the tetragonal bipyramid 4a with C<sub>s</sub> symmetry is found to be the most stable isomer, which is in agreement with the results for V<sub>2</sub>Si<sub>4</sub><sup>-</sup> [29]. The isomer 4b with C<sub>2</sub> symmetry is a triangular prism structure, and its energy is 0.20 eV higher than that of the 4a isomer. Other low-lying isomers (4c and 4d), which have the structure similar to the 4a isomer, are higher in total energy than the lowest-energy isomer 4a by 0.74 and 0.83 eV, respectively. At  $n = 5$ , the isomers 5a and 5d are two structures derived from the low-lying isomers of Mg<sub>2</sub>Si<sub>4</sub> clusters, where the fifth silicon or magnesium atom is added into the Mg<sub>2</sub>Si<sub>4</sub> clusters at different site. The optimized results indicate that the isomer 5a with C<sub>s</sub> symmetry is the ground state structure. When the 4b isomer is capped with one silicon atom, the metastable isomer 5b (C<sub>1</sub>) is generated, which has an energy higher than isomer 5a by 0.10 eV. The isomer 5c (C<sub>1</sub>) is a pentagonal bipyramid, and its energy is higher than that of the 5b isomer by 0.01 eV.

For Mg<sub>2</sub>Si<sub>6</sub> clusters, a lot of possible configurations and spin states are tried to identify the ground state. The isomer 6a (C<sub>2</sub>) is evaluated to be the ground state structure, which can be regarded as capping one silicon atom on the lowest-energy isomer 5a. The geometry of isomer 6b (C<sub>1</sub>) can be created from the metastable isomer 5c by capping one magnesium atom. The isomer 6c (C<sub>s</sub>) has a structure similar to the isomer 6a, the difference between them is the site of the magnesium atom. The metastable isomer 6d (C<sub>1</sub>) is formed by capping two silicon atoms on the metastable isomer 4d. The energy of 6a is lower than those of 6b, 6c, and 6d isomers by 0.48, 0.63, and 0.81 eV, respectively. According to the calculated total energies of the four Mg<sub>2</sub>Si<sub>7</sub> isomers, the isomer 7a with C<sub>1</sub> symmetry is proved to be the ground state structure among all the isomers. The metastable isomer 7b (C<sub>1</sub>) can be viewed as capping one atom on the metastable isomer 6b. The other isomers (Fig. 1 (7c,C<sub>1</sub>), (7d,C<sub>1</sub>)) have lower stability than the 7b isomer, and their energies are higher than that of the 7b isomer by 0.13 and 0.49 eV, respectively.

At  $n = 8$ , four kinds of structures can be verified to be the minima. The isomer 8a with C<sub>1</sub> symmetry is predicted to be the lowest-energy structure. The isomer 8b (C<sub>2</sub>) is obtained when the 7a isomer is capped with one silicon atom. And its energy is higher than that of the 8a isomer by 0.14 eV. Other isomers, such as Figure 1 (8c, C<sub>s</sub>) and (8d, C<sub>1</sub>) are 0.60 and 0.85 eV higher in energy than that of the ground state isomer 8a, respectively. The low energy isomers for Mg<sub>2</sub>Si<sub>9</sub> are plotted in Figure 1 (9a–9d). When one silicon atom is biccapped on the top of the pentagonal prism, the lowest-energy isomer 9a is obtained. The other low-lying isomers 9b, 9c, and 9d are 0.01, 0.10, and 0.88 eV higher in energy than that of the ground state, respectively. For  $n = 10$ , we found that isomer 10a which is generated after one silicon atom is capped on the top of lowest-energy isomer 9a, is optimized to be the ground state structure. The isomer 10b with C<sub>1</sub> symmetry is a 4-4-4 layer structure, which is expected to be the lowest-energy structure. However, its energy is higher than that of isomer 10a by 0.17 eV. The other isomers 10c (C<sub>1</sub>) and 10d (C<sub>1</sub>) have weaker stability than 10a because their energies are higher than that of the 10a isomer by 0.51 and 0.73 eV, respectively. Likewise, we identified four low energy isomers for Mg<sub>2</sub>Si<sub>11</sub>, as shown in Figure 1. The calculations turn out that the Si-centered cage-like geometry 11a is the lowest-energy structure in all isomers. The other low-lying isomers, including Figure 1 (11b, C<sub>1</sub>), (11c, C<sub>1</sub>), and (11d, C<sub>1</sub>), which are 0.04, 0.77, and 0.81 eV higher in energy than that of the 11a isomer, respectively.

According to the above results, the lowest-energy geometries of Mg<sub>2</sub>Si<sub>*n*</sub> clusters are found to adopt 3D forms for  $n = 3 - 11$  and 2D structures for  $n = 1 - 2$ . In addition, the Mg<sub>2</sub>Si<sub>*n*-1</sub> structure capped with one magnesium or silicon atom is the dominant growth pattern for Mg<sub>2</sub>Si<sub>*n*</sub> clusters of various sizes ( $n = 1 - 11$ ).

### 3.2. Relative Stabilities

To predict the relative stabilities of the lowest-energy Mg<sub>2</sub>Si<sub>*n*</sub> ( $n = 1 - 11$ ) clusters, we calculated the averaged binding energies [ $E_b(n)$ ], fragmentation energy [ $E_f(n)$ ], and second-order energy difference [ $\Delta_2 E(n)$ ] and summarized the theoretical results in Table 2. The  $E_b(n)$ ,  $E_f(n)$ , and  $\Delta_2 E(n)$  for Mg<sub>2</sub>Si<sub>*n*</sub> clusters can be expressed by the following formulas [30, 31]:

Table 2. Calculated averaged binding energies ( $E_b$ ), the second-order energy difference ( $\Delta_2E$ ), the fragmentation energy ( $E_f$ ), HOMO–LUMO energy gaps ( $E_{\text{gap}}$ ), and Frequencies of the lowest-energy structures for Mg<sub>2</sub>Si<sub>*n*</sub> ( $n = 1 - 11$ ) clusters.

<i>n</i>	$E_b$ (eV)	$\Delta_2E$ (eV)	$E_f$ (eV)	$E_{\text{gap}}$ (eV)	Frequency (cm <sup>-1</sup> )
1	0.53			1.51	120.1, 226.2, 293.3
2	1.19	-0.10	3.17	1.42	42.6, 126.2, 185.4, 193.1, 320.8, 516.3
3	1.79	-1.14	3.27	1.55	61.5, 110.4, 142.3, 198.2, 232.3, 303.1
4	2.07	3.33	4.41	2.37	70.4, 77.5, 158.0, 237.6, 279.7, 286.3
5	2.22	-0.74	3.08	1.75	72.4, 72.9, 134.9, 144.8, 149.2, 197.9
6	2.42	0.74	3.82	2.58	53.4, 55.6, 80.9, 124.2, 145.9, 165.6
7	2.49	-0.64	3.08	2.23	56.9, 83.0, 116.1, 135.1, 164.6, 167.8
8	2.61	0.68	3.72	2.56	65.2, 76.6, 95.8, 100.5, 125.1, 128.6
9	2.65	-0.13	3.05	2.43	58.2, 75.7, 93.4, 118.8, 124.4, 140.6
10	2.70	0.17	3.18	1.85	63.3, 84.1, 113.3, 116.9, 127.4, 139.2
11	2.72		3.00	1.84	68.4, 82.7, 106.4, 119.3, 123.4, 133.1

$$E_b(n) = \frac{nE_k(\text{Si}) + 2E_k(\text{Mg}) - E_k(\text{Mg}_2\text{Si}_n)}{n+2}, \quad (1)$$

$$E_f(n) = E_k(\text{Mg}_2\text{Si}_{n-1}) + E_k(\text{Si}) - E_k(\text{Mg}_2\text{Si}_n), \quad (2)$$

$$\Delta_2E(n) = E_k(\text{Mg}_2\text{Si}_{n+1}) + E_k(\text{Mg}_2\text{Si}_{n-1}) - 2E_k(\text{Mg}_2\text{Si}_n), \quad (3)$$

where  $E_k$  represents the total energy of the lowest-energy silicon atom, magnesium atom, Mg<sub>2</sub>Si<sub>*n*</sub> cluster, Mg<sub>2</sub>Si<sub>*n-1*</sub> cluster, and Mg<sub>2</sub>Si<sub>*n+1*</sub> cluster, respectively.

The  $E_b(n)$ ,  $E_f(n)$ , and  $\Delta_2E(n)$  values of the most stable Mg<sub>2</sub>Si<sub>*n*</sub> clusters at the B3LYP/6-311G (d) level of theory as a function of clusters size are depicted in Figure 2a and b. As shown in Figure 2, with the increasing size of clusters, the  $E_b(n)$  of Si<sub>*n+2*</sub> [20] and Mg<sub>2</sub>Si<sub>*n*</sub> clusters gradually increase. When  $n \leq 3$ , the  $E_b(n)$  rapidly increase as the cluster size of the pure Si<sub>*n+2*</sub> and Mg<sub>2</sub>Si<sub>*n*</sub> clusters increase, then the  $E_b(n)$  of the pure Si<sub>*n+2*</sub> and Mg<sub>2</sub>Si<sub>*n*</sub> clusters increases smoothly as the cluster size grows for  $n > 3$ . The results imply that Si<sub>*n+2*</sub> and Mg<sub>2</sub>Si<sub>*n*</sub> clusters continuously gain energy during the growth. The  $E_b(n)$  increases rapidly from Mg<sub>2</sub>Si (0.53 eV) to Mg<sub>2</sub>Si<sub>3</sub> (1.79 eV), corresponding to the structure transition from two to three dimensions. In addition, the  $E_b(n)$  of Mg<sub>2</sub>Si<sub>*n*</sub> doped counterparts is obviously smaller than that of the corresponding pure Si<sub>*n+2*</sub> [20] cluster, which hints that the impurity magnesium atoms can enhance the chemical activity of the pure silicon clusters. This phenomenon is similar to the  $E_b(n)$  of Cu<sub>2</sub>Si<sub>*n*</sub> clusters [32]. The fragmentation energy  $E_f(n)$  and second-order energy difference  $\Delta_2E(n)$  are sensitive indicators of the relative stability. Three prominent maxima  $E_f(n)$  for the most stable Mg<sub>2</sub>Si<sub>*n*</sub> clusters are found at  $n = 4, 6, \text{ and } 8$ , indicating that the lowest-energy Mg<sub>2</sub>Si<sub>4</sub>, Mg<sub>2</sub>Si<sub>6</sub>, and

Mg<sub>2</sub>Si<sub>8</sub> clusters have stronger relative stabilities compared to their corresponding neighbors. It is also found that the  $\Delta_2E(n)$  for the Mg<sub>2</sub>Si<sub>*n*</sub> clusters is bigger than other neighboring clusters when  $n = 2, 4, 6$  and  $8$ , indicating these clusters possess higher relative stabilities. This result is well consistent with the maxima presented in the fragmentation energy  $E_f(n)$ .

### 3.3. HOMO–LUMO Gaps and Charge Transfer

The energy gaps  $E_{\text{gap}}$  (The highest occupied-lowest unoccupied molecular orbital (HOMO–LUMO) energy gap) reflects the ability of electrons to jump from an occupied orbital to an unoccupied orbital, and represents the ability for the molecular to participate in the chemical reactions to some degree [33]. It is found that a larger gap associates with higher stability, in other words, a bigger gap signifies a weaker chemical activity. The  $E_{\text{gap}}$  for the lowest-energy geometry on the pure Si<sub>*n+2*</sub> [20] and doped Mg<sub>2</sub>Si<sub>*n*</sub> ( $n = 1 - 11$ ) clusters against the cluster size are plotted in Figure 2c. The data in Figure 2c show that the Mg<sub>2</sub>Si<sub>*n*</sub> clusters with  $n = 1, 4, 6, \text{ and } 9$  exhibit relatively wider  $E_{\text{gap}}$ , suggesting that these clusters are more stable than their neighbors, which is in agreement with the behavior of the fragment energy  $E_f(n)$  and second-order difference energy  $\Delta_2E(n)$  (except for  $n = 1$  and  $9$ ) shown in Figure 2a and b, respectively. So it can be concluded that the magic clusters are found at  $n = 4$  and  $6$  for the Mg<sub>2</sub>Si<sub>*n*</sub> clusters, reflecting that the Mg<sub>2</sub>Si<sub>4</sub> and Mg<sub>2</sub>Si<sub>6</sub> are the most stable geometries. Particularly, the Mg<sub>2</sub>Si<sub>6</sub> cluster has the largest  $E_{\text{gap}}$  of 2.58 eV compared with others, including the Mg<sub>2</sub>Si<sub>4</sub> cluster. Moreover, Figure 2c shows that the  $E_{\text{gap}}$  of the pure Si<sub>*n+2*</sub> clusters are larger than those of the Mg<sub>2</sub>Si<sub>*n*</sub> clus-

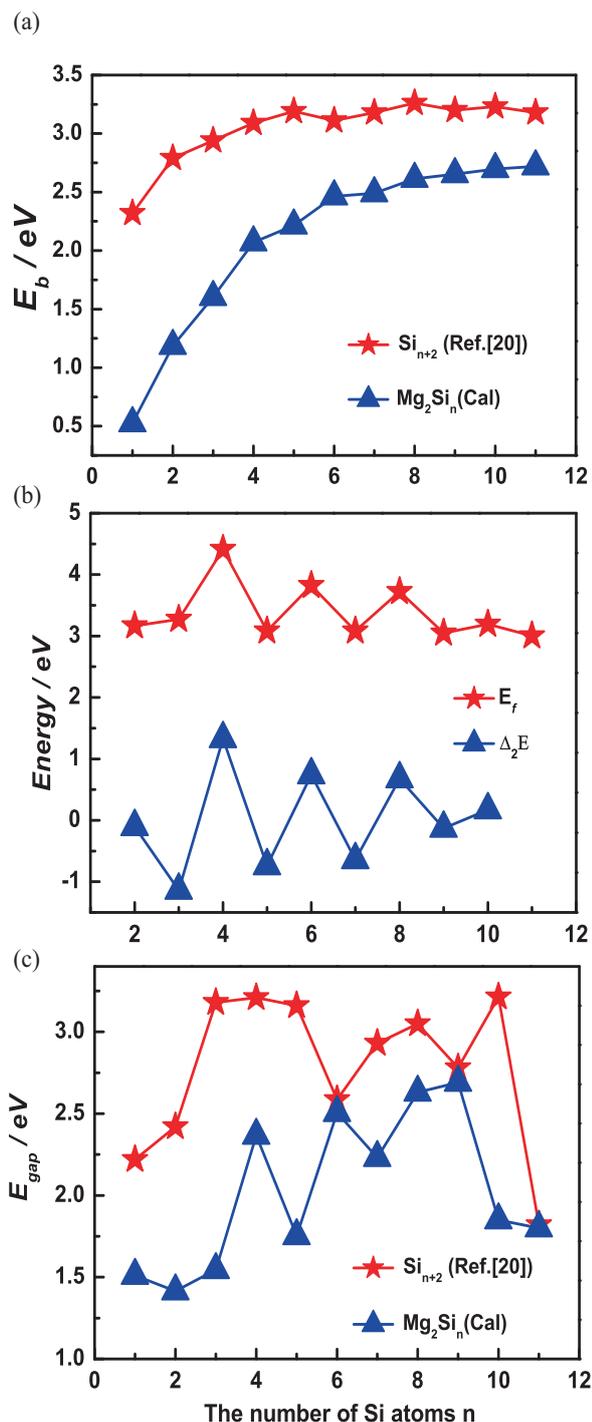


Fig. 2 (colour online). Size dependence of the averaged binding energy  $E_b$ , second-order difference energy  $\Delta_2 E$ , fragment energy  $E_f$ , and the HOMO–LUMO energy gap  $E_{gap}$  of the lowest-energy structures for Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1–11$ ) clusters.

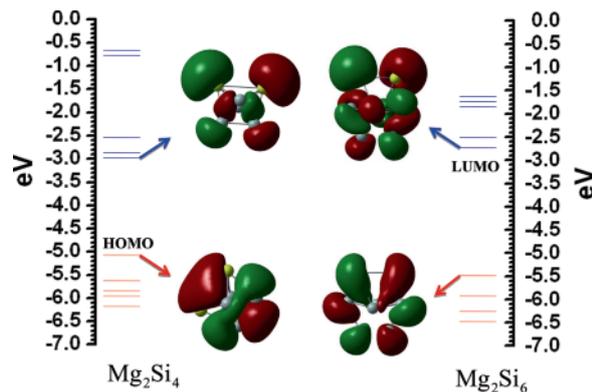


Fig. 3 (colour online). Calculated molecular orbital energy levels of Mg<sub>2</sub>Si<sub>4</sub> and Mg<sub>2</sub>Si<sub>6</sub> clusters together with the molecular orbital maps of the HOMOs and LUMOs.

ters, which means that the metallic characteristics of Mg<sub>2</sub>Si<sub>n</sub> clusters are enhanced by the doping of Mg atoms.

In order to investigate the effect of magnesium atoms doping on electronic structure of silicon cluster, the molecular orbital energy levels of the magic clusters for Mg<sub>2</sub>Si<sub>n</sub> clusters are calculated and displayed in Figure 3. The red line (lower half of the ordinate) shows the occupied orbital and the blue line (upper half of the ordinate) shows the unoccupied orbital. Compared with Mg<sub>2</sub>Si<sub>4</sub> and Mg<sub>2</sub>Si<sub>6</sub>, Mg<sub>2</sub>Si<sub>6</sub> is obviously characteristic of the molecular orbital producing the degeneration of the energy level in the vicinity of the HOMO–LUMO, and the energy level of the LUMO elevated strongly, which probably leads to its largest value of energy gap.

The natural population analysis (NPA) can provide reliable charge transfer information. So, the NPA analysis of magnesium atoms for the lowest-energy Mg<sub>2</sub>Si<sub>n</sub> clusters are calculated and listed in Table 3. From the table, one can see that the sign of the charges on magnesium atoms (0.388–0.861) is positive, while most of the silicon atoms possess negative charges. This result indicates that the charges transfer from the magnesium atoms to the Si<sub>n</sub> frame, that is, silicon atoms act as electron acceptor in all Mg<sub>2</sub>Si<sub>n</sub> clusters. This may be caused by reason that the electronegativity of magnesium (1.31) is much smaller than that of silicon (1.90) [34]; therefore, magnesium has a stronger ability to lose electrons.

Moreover, to understand the internal charge transfer, the natural electron configurations (NEC) of the mag-

Table 3. Natural population analysis (NPA) of magnesium atoms in the lowest-energy structures of Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters.

	1	2	3	4	5	6	7	8	9	10	11
Mg (1)	0.388	0.546	0.725	0.689	0.505	0.513	0.776	0.783	0.587	0.680	0.861
Mg (2)	0.388	0.372	0.725	0.689	0.505	0.513	0.776	0.783	0.405	0.843	0.747

Table 4. Charges of 3s, 3p, 4s, 3d, and 4p orbitals for the magnesium atoms in Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters.

		1	2	3	4	5	6	7	8	9	10	11	
Mg (1)	3s	1.34	0.75	0.86	0.85	1.01	1.14	0.62	0.87	1.14	0.52	0.66	
	3p	0.26	0.68	0.39	0.44	0.47	0.33	0.56	0.32	0.25	0.75	0.46	
	4s	0.01		0.01							0.01		
	3d	0.01	0.01	0.01	0.01	0.01	0.01	0.01	0.02	0.01	0.01	0.03	0.01
	4p	0.01	0.01	0.01	0.01	0.01	0.01	0.01	0.02	0.01	0.01	0.02	0.01
Mg (2)	3s	1.34	1.40	0.86	0.85	1.01	1.14	0.62	0.87	0.96	0.75	0.83	
	3p	0.26	0.22	0.39	0.44	0.47	0.33	0.56	0.32	0.61	0.44	0.40	
	4s	0.01		0.01						0.01	0.01	0.01	
	3d	0.01		0.01	0.01	0.01	0.01	0.01	0.02	0.01			
	4p	0.01		0.01	0.01	0.01	0.01	0.01	0.02	0.01			

nesium atoms in the lowest-energy Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters are calculated and summarized in Table 4. For a free magnesium atom, the configuration of valence electrons is 3s<sup>2</sup>. With regard to the impurities, the NEC values reveal that the 3s orbital loses 0.60–1.38 electrons, while the 3p orbital receives 0.22–0.75 electrons for Mg<sub>2</sub>Si<sub>n</sub> clusters. The contribution of the 4s, 3d, and 4p orbitals are nearly zero, so it can be neglected. Therefore, we can conclude that the charges in the lowest-energy Mg<sub>2</sub>Si<sub>n</sub> clusters transfer from the 3s orbital of the magnesium atom to the silicon atoms and 3p orbital of the magnesium atom. And the electronic charge distributions of Mg<sub>2</sub>Si<sub>n</sub> clusters are primarily governed by s- and p-orbital interactions. While the charges of 3s and 3p orbitals for two magnesium atoms in the lowest-energy Mg<sub>2</sub>Si<sub>n</sub> clusters are equal except for the Mg<sub>2</sub>Si<sub>2,9–11</sub> clusters.

#### 4. Conclusion

The equilibrium geometries, growth pattern behaviors, relative stabilities, and electronic properties of magnesium-doped silicon clusters have been obtained by using density functional theory at the B3LYP/6-

311G (d) level. In this work, we focused on analyzing a series of properties for these Mg<sub>2</sub>Si<sub>n</sub> clusters including averaged binding energies, fragmentation energy, second-order energy difference, and HOMO–LUMO energy gaps. The theoretical results show that the Mg<sub>2</sub>Si<sub>4</sub> and Mg<sub>2</sub>Si<sub>6</sub> are the most stable geometries in Mg<sub>2</sub>Si<sub>n</sub> ( $n = 1 - 11$ ) clusters. By analyzing the natural population and natural electronic configuration, it is found that the charges transfer from 3s orbital of the magnesium atom to the silicon atoms and 3p orbital of the magnesium atom, namely magnesium atoms act as electron donors in all Mg<sub>2</sub>Si<sub>n</sub> clusters.

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