

## Molecular Complexes, 13.

### Weak Arene Complexes of Nitroaromatics. $^1\text{H}$ NMR Studies of Conformations and Polar Effects. *syn-anti* Selectivity of 5-Nitrofurfural

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Z. Naturforsch. **57b**, 1305–1314 (2002); received July 23, 2002

Arene Complexes in  $\text{CCl}_4$ , Formation Constants, Topologies

The refined  $^1\text{H}$  NMR method (AUS concept,  $\text{CCl}_4$ ) provided association constants  $K$  and approximate complex shifts  $IK$  for complexes of nitroaromatics A with aromatic hydrocarbons D. 5-Nitrofurfural (**1**) and benzene (**B**) or toluene (**T**) form *syn-1B* or *syn-1T*, respectively; a proton dependence of  $K$  indicates about 10% *anti* complexes. The dimer  $\text{T}_2$  of **T** additionally yields some  $\text{1T}_2$  since  $IK$  values are 29–34% higher than for **1B**. With naphthalene (**N**) or phenanthrene all protons of **1** provide the same  $K$ . Picrylacetone (**2**) and **B** form **2B** and at least one **2BB** adduct; published parameters of **2B** and **2BB** are inconsistent. **2** and **N** or **M** (1,3,5-trimethylbenzene) form only 1:1 complexes. All results are in accord with  $\text{CH}_2\text{COMe}$  standing perpendicular to picryl. 2,4-Dinitrophenylacetone (**3**) in a rather flat conformation coordinates with a single stacking **N** centred over C-6; **3** yields two isomeric stacking complexes **3B** and at least one edge-on **3B** complex. 2-Chloro-5-nitropyridine (**5**) forms stacking complexes under repulsion of D by the ring N. In contrast to **1** and 4-nitrobenzaldehyde (**6**) 1,4-dinitrobenzene (**7**) and **T** form only a 1:1 complex that is stabilized by dipole-dipole attraction. Equal and unequal shielding of proton pairs (ortho or meta) in complexes of type **6B** is discussed.

## Introduction

Formation of complexes AD from compounds A and aromatic hydrocarbons D was often studied with NMR methods, in particular when A was a nitroaromatic compound [1–4]. The efforts concentrated on the determination of formation constants  $K$  but for weak complexes often with unsatisfactory results [2]. Consistent results for weak complexes and full exploitation of the potential inherent in  $^1\text{H}$  NMR can be reached when solvent effects are excluded as far as possible and when non-specific NMR effects resulting from the necessary excess of D are quantitatively considered both for A (see below) and for an internal reference [5, 4]. Disturbing solvent influences are avoided with the solvent  $\text{CCl}_4$  and non-linear effects on the reference signal are avoided by an external reference. Only one possible internal reference is known [6,7] whose *cis*-methyl signal position depends linearly on concentrations of D,

*viz.* 1-chloroisobutene (1-chloro-2-methylpropene) CLIB.

When the above requirements are taken into account (AUS concept) [5, 4], separate data reduction for each group of isochronic A protons provides  $K$  and an approximate complex shift  $IK$  arising from the complexing D by diamagnetic shielding of this proton group.  $K$  of weak complexes depends mainly on polar effects, dispersion forces and steric effects [1, 8]. With signals coming of different molecular sites of A their  $IK$  values usually contain informations about the mean topology of AD, primarily about the complex centre (mean point of A covered by the centre of D) in stacking complexes [1, 8]. Further influences on  $K$  and  $IK$  come of dipole-dipole and dipole-quadrupole interactions, impediment of torsional vibrations, formation of edge-on and of isomeric complexes [1]. Torsional vibrations in aromatic compound A make signals go a little upfield by decreasing the effects of resonance and magnetic anisotropies [1]. Distinct influ-

ences of conformations or of conformational flexibility in A were not observed so far. Some findings may be useful for studies of host-guest complexes.

## Method

A concise description of the AUS method was recently given [1]. This modification of the well-known Scatchard-Foster-Fyfe (ScFF) method arose from interrelated findings [6] with varying concentrations of benzene **B** ( $C_6D_6$ ) in  $CCl_4$ : (a) a non-linear ScFF plot turning from the negative to a positive slope for the least shifted signal of caffeine when internal TMS was the reference and (b) a quadratic term in the dependence of the TMS signal relative to an external signal. The very large concentration range possible with **B** (up to > 10 M) made such effects easily detectable. The small range with **D** = hexamethylbenzene was probably the main reason why the influence of internal TMS was overlooked in a previous comparison [9] of internal and external referencing. The reported shifts (INT) with internal referencing were always smaller and even after susceptibility corrections of the externally referenced shifts (EXT) there was a clear trend that was the more distinct the higher the concentration of this **D** was despite the small upper limit of 0.22836 mol/kg solution. INT was smaller than EXT by 0.1–1.2 Hz in 27 probes, equal to EXT in 4 probes and greater than EXT by 0.1 Hz in 3 probes.

The AUS concept takes into account non-specific shielding of A protons by D (total concentration  $[D_0]$ ) using linear shift corrections  $a_1 [D_0]$  for A and  $a_2 [D_0]$  for AD. The change of magnetic susceptibility with  $[D_0]$  is considered [4] by  $b [D_0]$ . Then, equations (1)–(5) replace the ScFF relation  $\Delta_0/[D_0] = -K \Delta_0 + K \Delta_{AD}$ . Definition of  $\Delta_0$  and  $\Delta_{AD}$  follows;  $Icpt$  replaces the intercept  $K \Delta_{AD}$ . Total A concentration  $[A_0]$  is made small while  $[D_0]$  in an experimental series of n solutions is evenly distributed over the largest possible range. A small  $[A_0]$  is advantageous for the methodical requirement  $[D_0] \gg [A_0]$  but it also makes chemical shifts less sensitive to possible self associations of A (*cf.* Table 3 in [1]) that may influence  $\Delta_0$  for low  $[D_0]$ .  $\Delta_0$  is the difference between chemical shifts of A in absence of D and in presence of  $[D_0]$ .  $IK$  is an approximation for  $\Delta_{AD} = (\delta_A \delta_{AD})$  of the classical model or, more precise [5,4],  $\Delta_{AD,00}$  of the AUS model.  $IK$  and  $m_2$

are obtained for each signal of A;  $m_1$  remains unknown but resembles  $m_2$  in size.  $IK$  depends on the diamagnetic anisotropy of D and on distances in AD enabling conclusions on the mean topology of AD. For further details and computer programs Sc-AUS, based on the ScFF method [10] and CA-AUS, based on the Cresswell-Allred method [11] see [1, 4, 9]. Sc-AUS is usually more sensitive to experimental scatter [12].

$$\begin{aligned} (\Delta_0 m_2 [D_0])/[D_0] &= K (\Delta_0 m_2 [D_0]) + Icpt & (1) \\ Icpt &= K [\Delta_{AD,00} + (m_1 - m_2) K^{-1}] & (2) \\ m_1 &= a_1 + b & (3) \\ m_2 &= a_2 + b & (4) \\ IK &= Icpt/K & (5) \end{aligned}$$

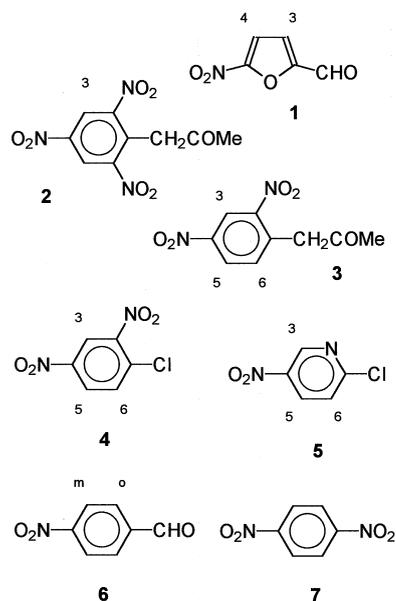


Fig. 1. Structures of compounds A; numbering and labels of proton sites.

## Results and Discussion

Results obtained with **1** (Fig. 1) are presented in Table 1, results with **2**, **3**, **5** and **7** in Table 2. **B** stands for deuterated benzene ( $C_6D_6$ ), **T** for deuterated toluene ( $C_7D_8$ ), **M** for mesitylene (1,3,5-trimethylbenzene), **N** for deuterated naphthalene ( $C_{10}D_8$ ) and **P** for phenanthrene. SF is the saturation fraction  $[AD]/[A_0]$ . Only CA-AUS results are reported, the Sc-AUS results did not significantly deviate.  $m_2$  values are reported but usually not discussed. Any temperature dependence of chemical shifts for CH signals is given in the discussion.

Table 1. Parameters of complexation of **1** from *m* experimental series; rounded parameters of complexation of **6**<sup>a</sup>.

AD	<i>m</i>	SF	<i>K</i> / M <sup>-1</sup> (max. dev.)	10 <sup>2</sup> <i>IK</i> (max. deviation) / ppm 10 <sup>2</sup> <i>m</i> <sub>2</sub> (max. deviation) / ppm M <sup>-1</sup>		
				CHO	3-H or <i>ortho</i>	4-H or <i>meta</i>
<b>1-B</b> total	1	0.11–0.62	0.149 (0.009)	140.6 5.9	147.4 (5.4) 8.0 (0.2)	144.8 (1.7) 7.4 (0.1)
CHO		0.140				
3-H + 4-H		0.154 (0.003)				
<b>1-T</b> total	2	0.10–0.54	0.138 ± 0.007	189.8 (1.0) 6.0 (0.2)	190.4 (0.6) 8.1 (0.1)	186.9 (0.7) 7.7 (0)
CHO		0.130 (0.009)				
3-H + 4-H		0.143 ± 0.001				
<b>1-N</b>	2	0.08–0.50	0.508 ± 0.011	108.5 (2.6) 20.1 (0.4)	152.7 (3.6) 26.2 (0.6)	144.7 (4.0) 22.9 (0.5)
<b>1-P</b>	3	0.11–0.44	0.814 ± 0.004	148.8 (2.2) 25.1 (1.4)	189.7 (1.4) 36.3 (1.3)	186.2 (5.8) 27.9 (2.3)
<b>6-B</b>			0.125	152	158	123
<b>6-T</b>			0.107	207	204	172
<b>6-N</b>			0.415	143	161	129

<sup>a</sup> Taken from [1].

### Complexes of 5-nitrofurfural (**1**)

Compound **1** is the furan analogue of 4-nitrobenzaldehyde (**6**) (Fig. 1) whose complexes were described recently [1]. The parameters (Table 1) of **1B** and **1T** show *K* from CHO to be 10% smaller than *K* from ring protons. A similar result with 1-chloro-2,4-dinitrobenzene (**4**) (Fig. 1) and **T** was found to come of isomeric complexes, *viz.* a stacking complex and a (minor) corner-on complex detectable by enlarged *IK* values [1] arising from corner-on complexes with **T** and its dimer **T**<sub>2</sub>. The 10% *K* difference for **1B** and **1T** result in another way from isomeric complexes.

Complexes of **1** should resemble those of **6** despite slightly different molecular geometries of 5- and 6-membered rings. In particular one may expect **1B** to have two characteristics of **6B** (Fig. 2: B), *viz.* [1] repulsion of **B** by the torsionally vibrating NO<sub>2</sub> and dipole-quadrupole interaction of **B** with H of the non-symmetrically vibrating CHO. But **1** can form isomeric complexes from its *syn-anti* rotamers (Fig. 2: A). A topology as shown in Fig. 2 (C) would suffer from repulsion of **B** both by the ring O and, less so, by NO<sub>2</sub> which is closer to **B** than in **6B**. Thus, **1** can only yield the complex

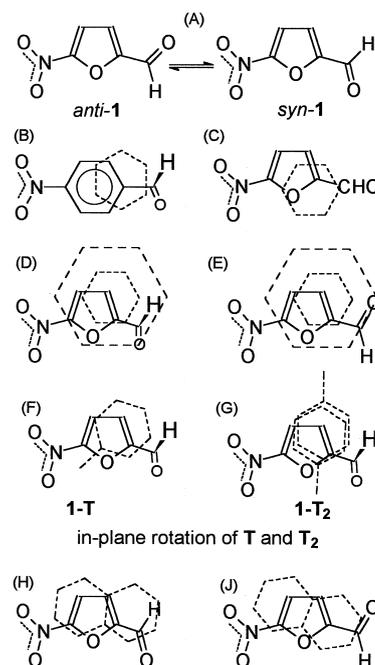


Fig. 2. (A) *syn-1* and *anti-1*, (B) **6B**, (C) unlikely complex centre of **1**, (D,E) complexing of **1** with **B** (approximate van der Waals dimensions), (F,G) **1T** and **1T**<sub>2</sub> (both shown for *syn-1* only), (H,J) **1N** (*syn* and *anti*).

topologies shown in Fig. 2 (D, E). The position of **B** required for a dipole-quadrupole interaction with H of CHO can only be reached with **syn-1B** (Fig. 2: D) making it the main complex. The corresponding topology (Fig. 2: E) of **anti-1B** is even destabilized by polar interaction of **B** with O of CHO. Diamagnetic shielding of the CHO proton by **B** is much smaller in this **anti-1B** than in **syn-1B**. Hence the small amounts of **anti-1B** will provide disproportionately small contributions to experimental  $\Delta_0$  values for CHO so that its  $K$  is caused nearly exclusively by the *syn* complex. There will be no substantial *syn-anti* difference (Fig. 2: D, E) in  $IK$  of the ring protons so that their  $K$  may be regarded as the sum of both  $K$  values. The same line of reasoning holds for the  $K$  difference obtained from **1** and **T**. No separate  $K$  values for the isomeric complexes can be obtained since each experimental signal is the weighted average of all dissolved species.

Apart from the *syn-anti* isomerism there is indeed a good analogy between complexes of **1** and of **6**. So, the  $IK$  values of **1T** are 29–34% higher than those of **1B** indicating complexing both with **T** and **T<sub>2</sub>** (Fig. 2: F, G) as reported [1] for **6**; the molecular dimensions (Fig. 2: F) exclude a dipole-dipole attraction between **T** and **1** that would make formation of **1T<sub>2</sub>** negligible. Further,  $K$  of the **T** complexes is smaller than  $K$  of the **B** complexes since repulsion of D by the vibrating NO<sub>2</sub> is increased by the in-plane rotation of **T** (and **T<sub>2</sub>**). Some differences between **1** and **6** remain. The higher  $K$  values for the complexes of **1** may arise both from less repulsion by the vibrating NO<sub>2</sub> owing to the molecular geometry and from stronger polar effects as shown by the higher partial charge (AM1) for the mainly covered HCCH moieties. Despite a likely upfield shift on going from **anti-1** to **syn-1** (*cf.* furfural below)  $10^2 IK$  for CHO of **1B** or **1T** is smaller (Table 1) than that one of **6B** or **6T**, respectively. This is in accord with small amounts of *anti* complexes and with a less covered CHO proton in *syn* complexes due to different molecular geometries. For  $IK$  values of ring protons see below. The temperature dependence for chemical shifts [1] of **1** confirms the mentioned torsional vibrations. Cooling a solution of **1** (0.016 M) in CCl<sub>4</sub> from 25 °C to –10 °C made the signals go downfield by 0.0165 ppm (CHO), 0.0501 ppm (3-H) and 0.0526 ppm (4-H) as com-

pared [1] to 0.0115, 0.040 (*ortho*) and 0.044 (meta) ppm for **6**.

No NMR data relevant to the *syn-anti* equilibrium constant are known for **1** in contrast to the parent furfural [13–15]. Dipole moments (in debye) of **1** are 2.60 (*anti*) and 6.07 (*syn*) from AM1 calculation, those of furfural [14] are 3.23 and 3.93 (microwave) or 2.67 and 3.67 (CNDO calculation). Ignoring any solvent interaction, AM1 predicts **anti-1** to be more stable by 5.5 kJ/mol while for furfural the *syn* isomer was more stable in dimethylether by 4.4 kJ/mol as calculated from the temperature dependence of coupling constants (4-H/CHO and 5-H/CHO) and hence of the *syn-anti* equilibrium [13]. Abraham and Siversns [14] have shown in a thorough review that the coupling 4-H/CHO in furfural (see below) decreases inversely with the polarity of the solvent in accord with the dipole moments. The analogous polarity influence on the *syn-anti* equilibrium of **1** can be neglected in the data reduction since *e.g.* for **1B** the permittivities of CCl<sub>4</sub> and **B** are 2.2 and 2.3. The *syn-anti* assignment of Dahlqvist and Forsén [13] for furfural has to be inverted [15, 14]. Then three final results [13] for furfural (dimethylether, –115 °C) are useful for a discussion of complexes of **1**. The spin coupling CHO/4-H is 0.85 Hz in the *anti* isomer but < 0.2 Hz in the *syn* isomer and the 3-H (CHO) signal of the *syn* isomer is 0.15 ppm downfield (0.18 ppm upfield) of the *anti* isomer. We now found for **1** in CCl<sub>4</sub> at 30 °C a 4-H/CHO coupling of 0.69 Hz indicating a predominance of **anti-1** under the assumption that it has about the same coupling constant (0.85 Hz) as *anti*-furfural. This equilibrium coupling decreases to 0.40 Hz when 20% of CCl<sub>4</sub> is replaced by the polar CDCl<sub>3</sub>. It decreases also on addition of **B**: to 0.62 Hz with [D<sub>0</sub>] = 4.8 M and to 0.55 Hz with [D<sub>0</sub>] = 9.0 M strongly supporting a preferred complexing of **syn-1**.

$K$  (Table 1) for **1N** or **1P** is independent of the proton studied; separate calculation for CHO and the ring protons provided  $K = 0.510$  and  $0.507 \text{ M}^{-1}$  for **1N** and  $K = 0.815$  or  $0.814 \text{ M}^{-1}$  for **1P**. This may come of either the absence of *anti* complexes or a decreased *syn* preference. The large arenes D may cover H of CHO so much better than **B** or **T** that the *syn* complexes, *e.g.* **syn-1N** (Fig. 2: H) are strong enough to prevent a detectable influence of *anti* complexes. Alternatively the polar interaction

Table 2. Parameters of complexation of **2**, **3**, **5** and **7** from *m* experimental series; rounded complex parameters<sup>a</sup> of **4B**.

A–D	<i>m</i>	SF	<i>K</i> / M <sup>-1</sup> (max. dev.)	10 <sup>2</sup> <i>IK</i> (max. deviation) / ppm 10 <sup>2</sup> <i>m</i> <sub>2</sub> (max. deviation) / ppm M <sup>-1</sup>				
				Me	CH <sub>2</sub>	3-H	5-H <sup>b</sup>	6-H
<b>2-B</b>	2	0.18–0.63	0.165	96.6 (2.0)				
			0.198	7.1 (0.03)	132.2 (1.6)			
		0.25–0.73	0.247		5.6 (0.09)	153.6 (0.5)		
<b>2-M</b>	2	0.12–0.48	0.489 ± 0.012		61.3 (1.3)	105.7 (0.8)		
					10.4 (0.5)	10.6 (0.02)		
<b>2-N</b>	2	0.15–0.55	1.206 ± 0.036	41.6 (0.7)	97.9 (2.3)	153.3 (1.3)		
				17.6 (0.6)	14.6 (0.4)	11.9 (0.17)		
<b>3-B</b>	2	0.19–0.51	0.107 (0.002)	113.2 (2.1)				
		0.19–0.59	0.134 ± 0.004	6.0 (0.1)	162.3 (6.4)	119.9 (1.9)	145.4 (6.1)	184.2 (3.6)
<b>4-B</b>			0.141			172	199	201
<b>3-N</b>	1	0.09–0.41	0.535 ± 0.020	63.8	149.1	98.0 (0.1)	153.6 (1.6)	200.9 (1.1)
				14.6	14.1	9.9 (0.01)	12.4 (0.22)	15.8 (0.17)
<b>5-M</b>	1	0.09–0.44	0.235 ± 0.002			96.7 (0.2)	174.6 (2.2)	159.6 (1.9)
						10.0 (0.1)	8.1 (0.2)	8.9 (0.1)
<b>5-N</b>	3	0.14–0.46	0.420 ± 0.003			93.6 (1.0)	201.4 (1.9)	192.0 (1.0)
						12.2 (0.1)	15.7 (0.4)	18.3 (0.6)
<b>7-T</b>	2	0.12–0.57	0.144 (0.06)					190.0 (2.0)
								6.4 (0.1)

<sup>a</sup> Taken from [1]. <sup>b</sup> One of the four lines could not be used with **3N**.

between a large D and CHO in the *syn* complex is rather offset by an increased contact interface in the *anti* complex (see *anti-1N* in Fig. 2: J) arising from the size of D and from hindered CHO vibrations. The low *IK* values of CHO are better compatible with the second possibility. Of course, one can expect rather low *IK* values owing to hindered vibrations but the two *IK* values for CHO are too small to arise exclusively in this way as shown by comparison with the ring *IK* values and as corroborated by *IK* comparison of **1N** with **6N** (Table 1). Moreover, the phenomenon of nearly equal *IK* values for **6B** and **6N** is also found with **1N** and **1B** except for the CHO *IK* of **1N** that is too low. *IK* values have the same order in **1P** as in **1N** and are greater in **1P** by 24–37% in congruence with the 30% difference for the above-plane ring current effects of both D [7].

### Complexes of (trinitrophenyl)- and (dinitrophenyl)acetone (**2** and **3**); complexes of 2-chloro-5-nitropyridine (**5**)

Complexing of picrylacetone (**2**) (Fig. 1) with **B** provided (Table 2) significantly differing *K* values for the three proton sites. This is incompatible with the pure 1:1 model. In addition to **2B** at least one complex **2BB** is formed whose **B** molecules are attached to different sites of **2**. Both sites may form isomeric 1:1 complexes of differing strength without an influence on *K*. Possible complex topologies including those for **2BB** are discussed in order to clarify why **2** and **N** do not yield **2NN** (see below).

Already a previous report [3] came to the conclusion that **2** forms complexes **2D** and **2DD** when **D** is benzene or a larger arene D; complex topologies and conformations of **2** were not discussed. Our parameters for **2** and **B** do not agree with

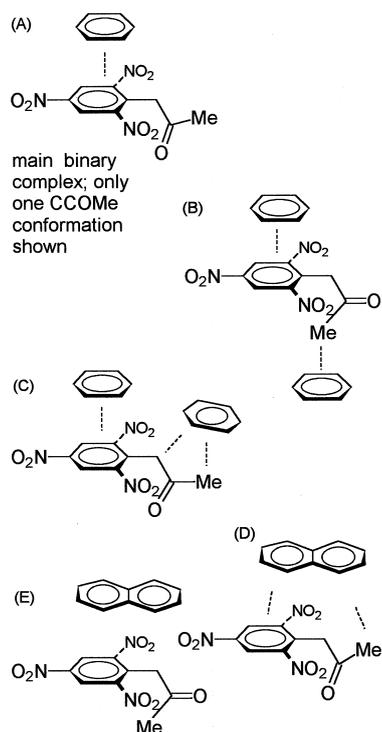


Fig. 3. Possible topologies for complexes of **2**.

those of this study, even with those calculated with the program Sc-AUS. The discrepancy is caused by at least three experimental details of the previous work. First, it is the internal reference (TMS) employed in contradiction to [6] and to the AUS concept as developed in [5]. Sc-AUS must not be used with internal TMS (see section Method). Second, it is the unusual concentration scale, *i.e.* mol/kg solution. There are good arguments [12] that in such studies molarity is the scale of choice. The unit “mol/kg solution” approximates molarity at very high concentrations if they can be reached. On the other hand  $K$  from mol/kg solution approximates  $K_{\text{molal}}$  for dilute solutions [16]. With weak complexes  $[D_0]$  has to cover the largest possible concentration range and then there is no linear relation between the scales; densities of **D** ( $< 1$  g/ml) and solvent (1.59 g/ml) differ too much. For  $\text{C}_6\text{H}_6$  in  $\text{CCl}_4$  1 M corresponds to 0.65, 2 M to 1.36 and 10 M to 10.43 mol/kg solution. Third, it is the use of  $\text{C}_6\text{H}_6$  that limited the concentration range of **D** to 0.3–1.5 mol/kg solution equal to 0.47–2.18 M as compared to 1.4–11 M for **2B** in Table 2. Nevertheless, apart from  $a_2$  values ( $a_2 =$

$m_2 - b$ , eq. (4)) some results show surprising similarities. As compared to the previously from Sc-AUS obtained  $K$  values the doubled  $K$  values of Table 2 are numerically equal ( $\text{CH}_2$ ), greater by 6% (Me) and smaller by 9% (3-H).  $IK$  values are given as  $\Delta_0$  in the previous report; they are smaller but have the same order  $3\text{-H} > \text{CH}_2 > \text{Me}$  as those of Table 2, *viz.* (given as  $10^2 IK$ ) 109, 108 and 89 ppm. However, the results of Table 2 show much more distinctly that the main centre for 1:1 and 1:2 complexes of **2** lies in the picryl moiety. One would expect that **2B** or most of its isomers have **B** stacking over picryl (Fig. 3: A), possibly shifted towards  $\text{CH}_2$  whose positive partial charge is increased by COMe.

Attempts to obtain reliable separate  $K$  values ( $K_1$  and  $K_2$ ) for steps **2 + D** and **2D + D** are doomed to failure since non-specific shielding by **D** must not be ignored. The previous attempt gave inconsistent results since protons of **2** were more (factor 1.3–2.4 for **D** =  $\text{C}_6\text{H}_6$ ) shielded by the second **D** than by the first **D**, *i.e.* each complex shift of **2DD** was more than twice that one of **2D** [3].

Intramolecular motions of previously studied compounds **A** at ordinary temperatures were rotations of methyl groups and torsional vibrations of other groups [1]. In this respect, **2** is a new type of **A** but its  $\text{CH}_2\text{COMe}$  cannot be flexible since steric requirements (confirmed by space-filling models) and polar effects make it stand perpendicular to the plane of the picryl moiety. Then the main binary complexes may be exemplified by Fig. 3 (A). The oxygen-oxygen repulsion between CO and  $\text{NO}_2$  favours the two conformations used for **2-BB** in Fig. 3 (B, C). They enable to attach one **B** to the sterically non-shielded face of picryl and to attach a second **B** over Me as shown in Fig. 3 (B, C) where the two **B** molecules of each ternary complex are sufficiently apart from one another even when van-der-Waals dimensions are considered. The second **B** in Fig. 3 (C) is drawn such as to form a T-shaped complex with  $\text{CH}_2\text{COMe}$  that will be stronger (*cf.* below the T-shaped **2B**) than the alternative complex with the second **B**. Thus, the association of **2** with **B** is likely to form only one ternary complex and two or three binary complexes, *i.e.* the picryl complex with one or two conformations of  $\text{CH}_2\text{COMe}$  plus a T-shaped complex over  $\text{CH}_2\text{COMe}$  arising from polar interactions

(CO dipole, CH<sub>2</sub>, Me). Apart from the latter complex the diamagnetic shielding of the picryl protons by **B** must always be strong probably making *IK* of 3-H a fair approximation for all other complexes. *K* provided by 3-H (0.24 M<sup>-1</sup>) might then possibly come close to *K*<sub>1</sub> that counts the first step (**2** + **B**).

**2N** (Table 2) provided the same large *K* from all proton sites showing the absence of a ternary complex. In the previous study of the naphthalene (C<sub>10</sub>H<sub>8</sub>) complex the pure 1:1 model was rejected despite finding *K* independent of the proton site [3]. The reported Sc-AUS parameters including *a*<sub>2</sub> values differed from our parameters due to the same kind of reasons as described above. *IK* values were smaller by 21–34% than those of Table 2; the doubled *K* of Table 2 numerically matches the previous result (2.3–2.4 kg/mol<sup>-1</sup>). The concentrations of D were in the range 0.03–0.30 mol/kg solution equal to 0.048–0.466 M as compared to 0.15–1.02 M for **2N** in Table 2. The 1:1 model was rejected because *a*<sub>2</sub> values did not have the expected [5] order Me > CH<sub>2</sub> > ArH. The wrong sequence is caused by non-specific shielding of internal TMS that makes the experimental Δ<sub>0</sub> unqualified for AUS corrections. Our *m*<sub>2</sub> values have the expected sequence; *a*<sub>2</sub> and *m*<sub>2</sub> differ only by the constant *b* (eq. 4). Pure 1:1 complexing with **N** in contrast to multi-complexing with **B** appears reasonable since van-der-Waals dimensions of **N** prevent the formation of a 1:2 complex of the type shown in Fig. 3 (C). Considering these dimensions and the dominance of stacking over picryl suggests that only two isomeric stacking complexes (Fig. 3: D, E) may be formed. A T-shaped complex over CH<sub>2</sub>COMe can scarcely compete with stacking whose contact interface with **N** is large making the stability contribution by dispersion forces great apart from a possible charge transfer contribution for stacking **2N**. *IK* values are also compatible with pure stacking, even *IK* of Me. **N** hinders the torsional vibrations of nitro groups thereby increasing the total magnetic anisotropy of picryl (*cf.* [1]). The conformation of **2** shown in Fig. 3 (E) puts Me into the shielding region of picryl. Stacking of **N** over the opposite face of picryl will increase the intramolecular shielding of Me in two ways. First, **N** + picryl shields more than **N** alone analogously to the strong shielding by **T**<sub>2</sub> in its complexes (*cf.* also the discussion of **3B**, below). Second, impeded vibrations increase the anisotropy of the picryl

moiety. The resulting small *IK* for Me is indistinguishable from a direct shielding by **N**. Shielding of CH<sub>2</sub> by **N** is counteracted by hindered vibrations of the neighbouring nitro groups making this *IK* rather small. Of course also *IK* for 3-H will include a counteraction of hindered vibration. This must be expected and is also supported by comparison with **2B** since *IK* values for 3-H are identical in analogy (Table 1) to the *IK* values of **6B** and **6N**. This supports the above interpretations of the 3-H parameters of **2B**.

For **2M** (Table 2) the Me signal of **2** was hidden under **M**. The other two signals were undisturbed and provided results that are compatible with the pure 1:1 model and with topologies analogous to Fig. 3 (D, E). *K* places **2M** between **2N** and binary **2B** as expected. *IK* comparison of **2M** with **2N** shows the weaker diamagnetic shielding by **M**, comparison with 3-H *IK* of **2B** indicates hindrance of torsional vibrations that comes of in-plane rotation of **M**. Edge-on attachment of a second in-plane rotating **M** analogous to Fig. 3 (C) is prevented by steric hindrance as with **2N**. The results with **2N** and **2M** leave no doubt that in contrast to the previous [3] conclusions phenanthrene and pyrene also form only 1:1 complexes resembling **2N** in topology. The pyrene complex is strong enough to make the disturbance by internal TMS insignificant and to give the expected order of *a*<sub>2</sub> values [3].

Removal of one ortho-nitro group from **2** reduces the molecular symmetry, increases the conformational flexibility of CH<sub>2</sub>COMe and weakens polar effects making complexes of **3** (Fig. 1,

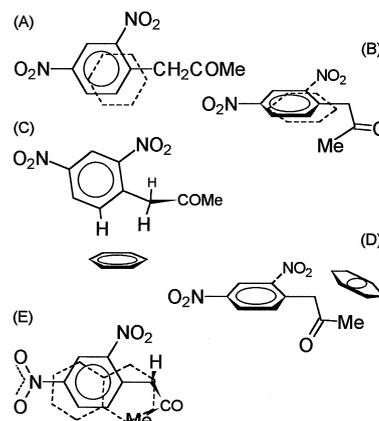


Fig. 4. (A, B) Stacking centre of **3**, (C, D, E) approximate complex topologies for **3**.

Table 2) weaker than those of **2**. The stacking centre of **2** is near the centre of its picryl moiety and is certainly equidistant to the ring protons, that one of **3** is moved away from the 1,4-diagonal (between carbons 1 and 4) of the phenyl ring (Fig. 4: A) as shown below. **3** easily takes on any CH<sub>2</sub>COME conformation as long as the oxygens of CO and NO<sub>2</sub> do not come close to each other. Complexing, however, may select a certain conformation so that polar effects and dispersion forces are optimized. This idea of a selected or perhaps even forced conformation seems to be novel in the field of weak complexes. There is no indication for a ternary complex with **3**.

Complex **3N** is discussed first (Fig. 4: E). The complex centre is near C-6 as shown by the *IK* sequence 6-H > 5-H > CH<sub>2</sub>. The last two have certainly been lowered by impeded torsional vibration of 4-NO<sub>2</sub> (5-H) and of 2-NO<sub>2</sub> (rather small influence on *IK* of CH<sub>2</sub>); this hindrance is probably more effective for 4-NO<sub>2</sub> than for 2-NO<sub>2</sub>. The low *IK* of 3-H is in accord with the relatively large distance to the complex centre and with impeded vibrations of both NO<sub>2</sub>. Fig. 4 (E) takes also two details into account. First, only one H of CH<sub>2</sub> is near **N** reducing the joint *IK* of CH<sub>2</sub>. Second, in this conformation Me is shielded by **N** but is put into the deshielding plane of dinitrophenyl making *IK* small. The joint polar effects of CH and Me are optimized in this conformation and even the CO dipole is favourably adjusted so that the stacking **N** simultaneously can complex with CHCOME in an edge-on manner.

**3B** (Table 2) yielded the same *K* from four of the five proton sites; the deviation for Me is moderate but requires an explanation. The largest *IK* of **3B** is that one of 6-H followed by CH<sub>2</sub> and 5-H again suggesting a stacking centre near 6-C although *IK* of CH<sub>2</sub> is too high as compared to **3N**. A marked polar interaction of **B** with CH<sub>2</sub> seems to be essential for this stacking centre since for the related compound **4** (Fig. 1) a rather centralized stacking has been reported [1] together with some edge-on complex standing with 5-H and 6-H on **B**. A similar isomeric edge-on complex can be expected for **3B** but **3** will stand with 6-H and one H of CH<sub>2</sub> on **B** (Fig. 4: C). This edge-on **3B** is probably stronger than that one of **4** (dipole moment 3.0–3.31 debye) [1] as follows from the dipole moment of 2,4-dinitrotoluene (4.33–4.38 debye) [17]. The small *K* for this edge-

on **3B** adds to *K* of the stacking complex for all proton sites except for Me providing a smaller total *K* for Me as found. This appears more reasonable than a ternary complex which influences only one of the five signals. With **3N** *K*<sub>stacking</sub> is strong enough to make *K*<sub>edge-on</sub> insignificant (*cf.* **4B** and **4T** in [1]). The total *K*s of **3B** and of **4B** (Table 2) are well comparable while, ignoring CH<sub>2</sub>COME, the *IK* values of **3B** are lower than those of **4B**, distinctly lower for 3-H due to the centralized topology [1] of the stacking **4B**. Contributions to *IK* by edge-on **3B** primarily depend on the distance between the proton and the centre of the complexing **B**. The resulting shielding by **B** is highest for 6-H and one H of CH<sub>2</sub>. It decreases in the sequence 5-H >> 3-H > Me. The stacking **B** is too small to cover Me analogously to **3N** (Fig. 4: E) so that this direct shielding of Me by **B** in both isomeric **3B** is small in contrast to the mediocre *IK* in Table 2. But both **3B** may have Me put rather beneath dinitrophenyl as shown for stacking **3B** in Fig. 4 (B) resulting in some intramolecular shielding for both. This shielding by picryl in stacking **3B** would be increased by **B** (two stacking rings in analogy to shielding by **T**<sub>2</sub>; *cf.* **2N**). In stacking **3B** this conformation enables an optimal polar interaction of CH<sub>2</sub> with 2-NO<sub>2</sub>. Part of the relatively high *IK* of Me (and partially of CH<sub>2</sub>) may also come of an edge-on complex over CH<sub>2</sub>COME (Fig. 4: D). A ternary stacking plus edge-on complex of type **2BB** is prevented by **B** of stacking **3B** due to the complex centre (Fig. 4: A).

The centralized stacking for complexes of **4** probably results from repulsion of D by the four O and from attraction of D by the two N and the proton between them. In chloronitropyridine (**5**) (Fig. 1) **5** (Table 2) one CNO<sub>2</sub> of **4** with its peripheral negative charges is replaced by a ring N providing a much more central negative charge. This shifts the complex centre near C-5, *i.e.* most remote from the ring N both for **5M** (Fig. 5: A) and **5N** (Fig. 5: B) as is evident from the *IK* values of both complexes even if *IK* for 6-H of **5M** might be less reliable than usual owing to partial signal overlap of 6-H and **M**. *IK* values are greatest for 5-H and smallest for 3-H (next to ring N). The latter *IK* is nearly equal in both complexes due to hindrance of NO<sub>2</sub> vibration by **N** and lack of hindrance in **5M** since in-plane rotation of **M** will be prevented by polar interactions of two methyl groups with negative partial charges in **5** (Fig. 5:

A). The  $IK$  difference between 6-H and 5-H coming of the distance from the centre of **D** is a bit greater in **5M** than in **5N** since **N** covers 6-H better than **M** does (Fig. 5: A,B).

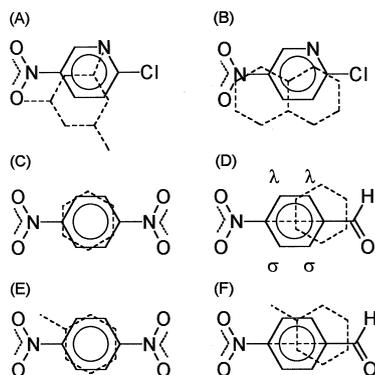


Fig. 5. (A, B) Complexes **5M** and **5N**, (C) complex centre of **7**, (D) complex centre of **6** and unequal shielding of proton pairs in **6**, (E,F) molecular dimensions for dipole-dipole interactions of **T** with **7** and **6**.

#### ***Diamagnetic shielding in complexes of 1,4-dinitrobenzene (7) and of 4-nitrobenzaldehyde (6); unequal shielding in isochronous proton pairs***

The highest  $10^2 IK$  (186 ppm) [4] of a stacking **B** complex was found with **7** (Fig. 1). This value is distinctly higher than  $10^2 IK$  values of **6B** both for meta and for ortho protons (Table 1). In contrast to **6** (Fig. 5: D) [1] the complex centre of **7** should be the centre of this molecule (Fig. 5: C). This can be confirmed as follows. **6** associates both with **T** and **T<sub>2</sub>** (high  $IK$ , see above) but **7** cannot form **7T<sub>2</sub>** since the postulated complex centre enables a dipole-dipole attraction (Fig. 5: E) that stabilizes **7T** and destabilizes **7T<sub>2</sub>** (stacking of three dipoles, *cf.* [1]); no such attraction is possible in **6T** (Fig. 5: F). Indeed,  $IK$  of **7T** (Table 2) is insignificantly higher than  $IK$  of **7B**. The difference of 2% is near the methodical limit but a small increase of diamagnetic shielding may even be expected from the diamagnetic anisotropies [7] of **T** and **B**. Dipole-dipole attraction makes **7T** a stronger complex than **7B** ( $K = 0.130 \text{ M}^{-1}$ ) [4] while **6T** is weaker than **6B** (Table 1) since the repulsive action of the vibrating  $\text{NO}_2$  is increased by Me of in-plane rotating **T**. This rotation is prevented in **7T**. A different torsional behaviour of  $\text{NO}_2$  groups in free **6** and **7**

can be excluded since the signals of neighbouring protons go downfield on cooling ( $\text{CCl}_4$ ,  $25^\circ\text{C} \rightarrow -15^\circ\text{C}$ ) by 0.043 ppm for **7** and [1] 0.044 ppm for **6**.

All protons of **7** are isochronous both in free **7** and in **7B**; each proton provides the same contribution to  $IK$ . **B** in **6B** is closer to CHO than to  $\text{NO}_2$  making ortho  $IK$  greater than meta  $IK$ ; but **B** is also moved away from the 1,4-diagonal of the benzene hexagon making shielding of each proton pair unequal (Fig. 5: D). A  $\sigma$  proton is more remote from the centre of **B** than the  $\lambda$  proton. The  $\sigma$  proton contributes to the joint  $IK$  less than the  $\lambda$  proton does. Shielding of the  $\lambda$  proton and hence its large contribution to the joint  $IK$  will not change much relative to a topology that has **B** on the 1,4-diagonal. The total result would be a comparatively small joint  $IK$ . Thus, **6B** and **7B** suggest that  $IK$  of such proton pairs decreases when the stacking centre is off the 1,4-diagonal. This idea is corroborated by complexes of **2** and **3**. The complex centre is on the 1,4-diagonal for the former and distinctly away from it for the latter (see above).  $10^2 IK$  for 3-H and 5-H of **2N** (both 153 ppm) is split into 98 and 184 ppm with **3N** averaging to 141 ppm, *i.e.* less than  $10^2 IK$  for **2N** although hindrance of  $\text{NO}_2$  vibrations is stronger in **2N** than in **3N**. The same comparison for the **B** complexes gives 154 (**2**) and 133 ppm (average for **3**) but is less reliable since 1:1 stacking is not the only kind of complexing both for **2B** and **3B**.

A  $\sigma$  contribution should be distinctly smaller than  $IK$  from **B** on the 1,4-diagonal according to the ring current model of Johnson and Bovey [18]: diamagnetic shielding decreases markedly even more than 35 nm above the plane of benzene when the distance to the central axis exceeds the ring radius. On going from 1 to 1.5 ring radii 36 nm above the benzene plane the shielding decreases by 0.27 ppm according to this model; then the contributions to  $10^2 IK$  for a  $\lambda$ - $\sigma$  pair may easily differ by about 20–30 ppm. At more central positions the shielding difference is smaller, *e.g.* 0.09 ppm on going from 0 to 0.5 ring radii so that the two contributions to  $10^2 IK$  of a proton pair insignificantly differ; the difference is zero when the centre of **B** is exactly over the 1,4-diagonal. The non-symmetrical vibration of CHO may furthermore decrease the contribution by the  $\sigma$  proton of the ortho pair in **6B**. It is essential to keep in mind that this reasoning concerns  $IK$  values, *i.e.* the changes of shifts caused by complexing.

The inherent shift difference of a proton pair in a frozen planar conformation of **6** should be equal for free **6** and for **6B**.

The furan analogue **1** of **6** has no  $\sigma$  protons. Complexes of **1** with **B**, **T** and **N** (Table 1) in contrast to those of **6** have nearly no difference between *IK* values of the ring protons although the complex centres of **1** and **6** (Fig. 5: D) are comparable. From the presence of the proton pairs in **6** one expects *IK* for ring protons to be greater in complexes of **1**. This is observed (Table 1) for protons next to NO<sub>2</sub>, *i.e.* 4-H of **1** vs. *meta* protons of **6**. *IK* values for protons next to CHO are even a bit smaller for **1** complexes. Here the effect of unequal shielding in complexes of **6** is overcompensated in complexes of **1** by the different *syn-anti* ratios of free and complexed **1**. Free **1** (mainly *anti-1*) will have the 3-H signal at a relatively high field (in *anti*-furfural 0.15 ppm upfield of *syn*-furfural, see above) making the shift difference between total free **1** and total complexed **1** (mainly *syn-1*) small. This may be the reason why *IK* values for ring protons next to CHO do not differ much in complexes of **1** and **6** with the same D.

### Experimental Section

For procedures, instrument, external reference, further details and computations see [1, 4, 6, 8]. The accuracy of  $\Delta_0$  was  $\leq 0.0016$  ppm; temper-

ature was  $30.0 \pm 0.3$  °C;  $[A_0]$  was 0.032–0.036 M for **1P**, 0.0632 M for **5M** and 0.0105–0.0222 M for the other complexes. The preparation of n (11–18, 23 for **1B**) solutions was conducted with the utmost accuracy from two stock solutions of known density at  $20 \pm 0.5$  °C by means of a microsyringe under mass control and avoidance of volatilization losses. The same capillary for the external reference (dioxane with **M**) was used throughout a series. Chemicals were used as purchased; **2** and **3** were prepared following described syntheses [3, 19]. The precision of the temperature dependence of shifts (Bruker Avance 300 instrument, internal TMS) was limited by changes of the respective multiplet.

Chemical shifts  $\Delta_0$  were measured in Hz and used as such in all calculations. Hz was converted to ppm in the final results. Each  $\Delta_0$  was the mean of three measurements unless the first two coincided. Shifts  $\Delta_0$  of doublets were taken from the stronger line unless the weaker line had a lower least square sum (SDDQ in the programs);  $\Delta_0$  of all lines were used for **1B** and for complexes of **3** and **5**. *K* is the mean of all lines used as described above; *IK* and  $m_2$  are means of the separately found values.

When cooled to  $-10$  °C a 0.15 M solution of **1** in CCl<sub>4</sub> becomes cloudy and shows a second CHO signal 0.40 ppm upfield (new ring signals < 0.1 ppm upfield). Both findings may possibly arise from secretion and self-association of highly polar *syn-1*.

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