

# The Structure of Gallium Phosphate Glasses by High-energy X-ray Diffraction

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X-ray diffraction experiments are used to obtain short-range order information of gallium phosphate glasses of meta- and pyrophosphate compositions. Parameters of the first-neighbor peaks, such as coordination numbers and distances, are obtained. A strong decrease of the Ga-O coordination number from  $6.0 \pm 0.2$  to  $4.6 \pm 0.2$  upon  $\text{Ga}_2\text{O}_3$  addition is found, which is accompanied by a shortening of the Ga-O distances. Only  $\text{GaO}_6$  octahedra exist at the metaphosphate composition. Close to the pyrophosphate composition, the majority of Ga atoms occupies already tetrahedral sites. The Ga-O coordination number behaves equivalent with the ratio  $M_{\text{TO}} = n(\text{O}_\text{T})/n(\text{Ga})$ , thus, with the number  $n(\text{O}_\text{T})$  of terminal oxygen atoms ( $\text{O}_\text{T}$ ) in phosphorus- $\text{O}_\text{T}$  bonds which are available for the coordination of each Ga atom. Thus, P- $\text{O}_\text{T}$ -Ga bridges are formed for all  $\text{O}_\text{T}$  atoms. The  $\text{GaO}_n$  polyhedra neither share  $\text{O}_\text{T}$  atoms nor form Ga-O-Ga bridges. With increasing fraction of  $\text{GaO}_4$  tetrahedra and decreasing lengths of the phosphate chains the network expands.

*Key words:* X-ray Scattering, Short-range Order, Phosphate Glasses.

## 1. Introduction

The oxides of gallium and phosphorus,  $\text{Ga}_2\text{O}_3$  and  $\text{P}_2\text{O}_5$ , are known as network-forming materials [1] where  $\text{GaO}_4$  and  $\text{PO}_4$  tetrahedra are the basic structural units. The two oxides can complement each other. Phosphorus can donate electron charge to gallium which would stabilize tetrahedral environments of both atoms. This mechanism is well developed in  $\text{GaPO}_4$  crystals [2, 3] where equal fractions of oxides are mixed and structures similar to silica polymorphs are formed.

The glass-forming range of the  $\text{Ga}_2\text{O}_3$ - $\text{P}_2\text{O}_5$  system is, however, limited to compositions rich in  $\text{P}_2\text{O}_5$  [4]. Following the rules of network changes for phosphate glasses [5] addition of modifying oxide leads to the disruption of P-O-P bridges where, at first, three-connected  $\text{PO}_4$  branching units are transformed to two-connected middle units before single-connected  $\text{PO}_4$  end units or, in the end, isolated  $\text{PO}_4$  units occur. At metaphosphate composition, i. e. in

glasses of 25 mole%  $\text{Me}_2\text{O}_3$  (with the more general formula  $\text{Me}_{1/v}\text{PO}_3$ , where  $v$  is the valency of Me), chain and/or ring structures of  $\text{PO}_4$  middle groups are formed. With further modifier addition the chains become shorter and at the pyrophosphate composition, i. e. in glasses of 40 mole%  $\text{Me}_2\text{O}_3$  (with the more general formula  $\text{Me}_{4/v}\text{P}_2\text{O}_7$ ) the  $\text{PO}_4$  dimers dominate. The corresponding changes are widely proved by results of  $^{31}\text{P}$  magic angle spinning nuclear magnetic resonance (MAS NMR) and vibrational spectroscopy. Many of such results are reviewed in [6].

The only known crystal structure in the range of glass formation with Ga atoms as the third species is that of  $\text{Ga}(\text{PO}_3)_3$  polyphosphate [7], where infinitely long, meandering phosphate chains are interconnected via  $\text{GaO}_6$  octahedra. A similar structure is also known of an  $\text{Al}(\text{PO}_3)_3$  modification [8]. Another detail of these structures is the fact that all oxygen atoms are either in P-O-P or P-O-Ga(Al) bridges where, consequently, the  $\text{GaO}_6$  ( $\text{AlO}_6$ ) octahedra

don't share any oxygen. Obviously, it is profitable for terminal P-O<sub>T</sub> bonds (O<sub>T</sub> = terminal oxygen) to share the double-bond character and to coordinate a second atomic neighbor. This principle requires a six-fold coordination of Ga(Al) atoms which is not common of Ga(Al) atoms in gallate(aluminate) glasses. Except of the P<sub>2</sub>O<sub>5</sub> polymorphs, no phosphate crystal is known which would possess 'dangling' P-O<sub>T</sub> bonds [9]. Obviously, P-O<sub>T</sub> bonds without a second neighbor of O<sub>T</sub> are reactive and, thus, are avoided. From this finding the same behavior was concluded to exist in phosphate glasses. This trend, and also the attempt to stabilize Me-O<sub>T</sub>-P bridging positions in symmetric MeO<sub>n</sub> environments were used to explain the density anomalies in phosphate glasses of bivalent modifier (Me) atoms [10, 11].

Rare earth oxides cause similar effects in binary phosphate glasses [12 - 14], which could be studied extensively only after applying a melt technique which avoids loss of the volatile P<sub>2</sub>O<sub>5</sub> component at the required high melting temperatures. The same problem limits the sample preparation also for Ga phosphate glasses, where the batches for ultra-phosphate samples which were melted in open crucibles result in metaphosphate glasses [15]. On the side of high Ga<sub>2</sub>O<sub>3</sub> fractions, Ga phosphate glasses are obtained up to the pyrophosphate composition [4, 15]. Changes of the PO<sub>4</sub> units in the corresponding glasses were analysed by vibrational spectroscopy [15], where also comparisons with the Ga(PO<sub>3</sub>)<sub>3</sub> crystal were made. If compared with Al phosphate glasses [16], the larger range of glass formation makes the Ga<sub>2</sub>O<sub>3</sub>-P<sub>2</sub>O<sub>5</sub> system more suitable to study structural changes in modifier environments. Ga atoms are good scatterers of X-rays. Ga-O bonds are longer than Al-O bonds, and thus are well resolved from P-O correlations. So X-ray diffraction experiments using high-energy photons are excellent in a search for changes in Ga-O environments. Though Al phosphate glasses were prepared [16] only in a small range with molar fractions,  $x$ , of Al<sub>2</sub>O<sub>3</sub> within  $0.31 < x < 0.35$ , significant changes of the Al environments were detected by <sup>27</sup>Al MAS NMR. With increasing  $x$  the fraction of AlO<sub>4</sub> tetrahedra increases at the expense of AlO<sub>6</sub> octahedra, whereby the clear fraction of AlO<sub>5</sub> units is about constant. In this work X-ray diffraction experiments are made to extract the short-range order information of Ga<sub>2</sub>O<sub>3</sub>-P<sub>2</sub>O<sub>5</sub> glasses in the range  $0.25 \leq x \leq 0.40$ , where the same sample material as prepared and studied in [15] is used.

## 2. Experimental

### 2.1. Sample Preparation

Details of the sample preparation are described in [15]. Batches of Ga<sub>2</sub>O<sub>3</sub> and NH<sub>4</sub>H<sub>2</sub>PO<sub>4</sub> were heated at 400°C and then melted in Pt crucibles for 20 min at 1300 - 1550 °C. Electron probe microanalysis revealed that, starting with batches containing  $0.15 \leq x \leq 0.30$  mole fractions of Ga<sub>2</sub>O<sub>3</sub>, glass samples close to metaphosphate composition ( $x = 0.25$ ) were obtained, while batches of 0.35 and 0.40 mole fractions Ga<sub>2</sub>O<sub>3</sub> yielded samples close to pyrophosphate composition ( $x = 0.40$ ). IR and Raman spectra confirmed that all glasses could be divided into the two groups [15]. One sample of each group was chosen for the X-ray diffraction experiments. After first Gaussian fitting using  $x = 0.25$  and 0.40 for the two samples, compositions have been changed to  $x = 0.27$  and 0.38 (cf. Chapter 3).

### 2.2. Diffraction Experiments

The X-ray diffraction experiments were performed at the BW5 wiggler beamline at DORIS III of Deutsches Elektronen-Synchrotron (Hamburg). An incident photon energy of 120 keV ( $\lambda = 0.0104$  nm) was chosen for the experiments. The beam size was  $1 \times 4$  mm<sup>2</sup>. Since the diameters of 2.5 mm of the silica capillaries (with wall thicknesses of 0.01 mm) containing the glassy powder are larger than the beam width exact absorption corrections are difficult. The scattering angles are small ( $2\theta = 28^\circ$  for  $Q_{\max} = 300$  nm<sup>-1</sup>) and the transmission coefficients are higher than 0.9, so that absorption is independent of the scattering angle.  $Q$  is the magnitude of a scattering vector with  $Q = 4\pi/\lambda \sin \theta$ . The electronic energy window of the solid-state Ge-detector was chosen to pass the elastic line and the full Compton peak but no fluorescence radiation. Dead-time corrections are made [17] with the parameter  $\theta = 1.08$  μs. A fraction of 0.91 of the incident photons is polarized horizontally. Corrections are made for background, container scattering, polarization and absorption. The scattering intensities are normalized to the structure-independent scattering functions which are calculated by a polynomial approach [18] of tabulated atomic data of the elastic scattering factors [19]. Compton scattering is calculated according to [20]. Finally, the Compton fraction is subtracted and the Faber-Ziman structure factors,  $S(Q)$ , are obtained [21, 22].

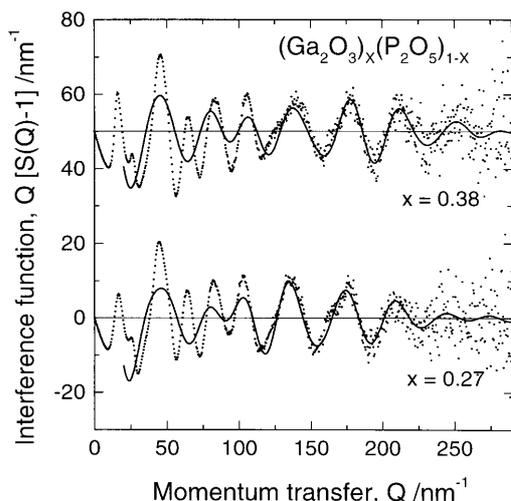


Fig. 1. Weighted interference functions,  $Q[S(Q)-1]$ , of the Ga phosphate glasses studied (experimental points: dots; model function calculated according to the first-neighbor peaks: solid lines).

### 3. Results

Weighted interference functions of the two samples are shown in Figure 1. Due to the use of high-energy photons (120 keV) of an intense source data are available up to more than  $300 \text{ nm}^{-1}$ , but the structure factors,  $S(Q)$ , showed reasonable behavior only up to  $\sim 270 \text{ nm}^{-1}$ . Data of higher  $Q$  are very noisy.

The correlation functions,  $T(r)$ , are obtained by Fourier transformation (FT) with

$$T(r) = 4\pi r \rho_0 + \frac{2}{\pi} \int_0^{Q_{\max}} Q[S(Q)-1] \sin(Qr) dQ, \quad (1)$$

where  $\rho_0$  is the number density of atoms. The values of  $\rho_0$  are estimated when performing the FT procedures. The  $r$ -range of the  $T(r)$  function in front of the first distance peak must oscillate about the zero line. Densities of  $70$  and  $80 \text{ nm}^{-3}$  were found to comply with this requirement for the glasses with  $x = 0.38$  and  $0.27$ , respectively. Note that the corresponding number densities of the related crystals  $\text{GaPO}_4$  [2] and  $\text{Ga}(\text{PO}_3)_3$  [7] are  $72.0 \text{ nm}^{-3}$  and  $85.2 \text{ nm}^{-3}$ , respectively. The final  $T(r)$  functions are shown in Fig. 2 where the  $Q_{\max}$  used in FT is  $270 \text{ nm}^{-1}$ .

Parameters of the first-neighbor distance peaks are determined by Gaussian fitting. The effects of the upper integration limit at  $Q_{\max}$  in the FT are taken into account [23]. Here, in case of the X-ray data the

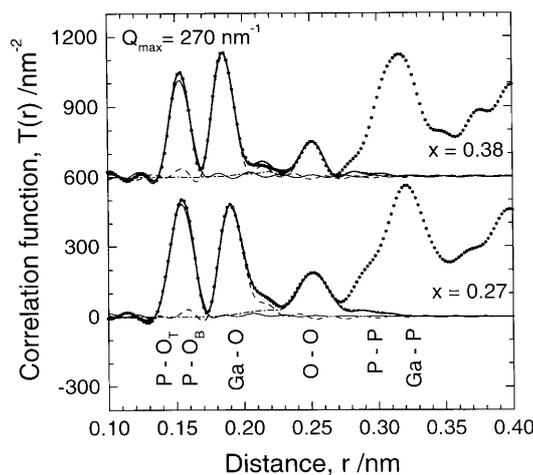


Fig. 2. Real-space correlation function  $T(r)$  of the Ga phosphate glasses studied (experimental data: dotted lines; model function obtained by Gaussian fitting: heavy solid lines). The partial contributions are given with P-O: thin solid lines; Ga-O: dashed lines; O-O: dash-dotted lines.

weighting factors depend on  $Q$ . For fitting the  $T(r)$  data, the Marquardt algorithm [24] was applied. Coordination numbers  $N_{ij}$ , mean distances  $r_{ij}$ , and full widths at half maximum (fwhm)  $\Delta r_{ij}$ , are the parameters of the model Gaussian functions. Two Gaussian functions are used for fitting the P-O peak with the  $\text{P-O}_T$  and  $\text{P-O}_B$  contributions ( $\text{O}_B$ : bridging oxygen atom). Split P-O peaks were found for all phosphate glasses which have been studied by neutron diffraction experiments of high resolving power ( $Q_{\max} \cong 500 \text{ nm}^{-1}$ ) [25]. The numeral ratio of P-O bonds depends on the degree of network depolymerization. The ratio is calculated from the compositions and is fixed in the fits. A possible asymmetry of the Ga-O peak is approximated by use of two Gaussian functions where the resulting two  $r_{\text{GaO}}$  have no special meaning and cannot be related to definite coordination polyhedra of Ga atoms. The O-O peak at  $\sim 0.25 \text{ nm}$  appears well resolved and free of other correlations. However, satellite ripples of the P-P and Ga-P correlations at  $0.30$  and  $0.32 \text{ nm}$  might somewhat affect the  $T(r)$  function in the  $r$ -range of the O-O peak. The resulting parameters of the first-neighbor model peaks are given in Table 1. Interference functions are calculated using these parameters and are compared with the experimental functions (cf. Fig. 1). A good agreement in the high- $Q$  range, where only short-range order information is reflected, can be stated. In Fig. 2

Table 1. Parameters resulting from Gaussian fitting of the first-neighbor P-O, Ga-O and O-O peaks of the correlation function,  $T(r)$ , which is shown in Figure 2.

Atom pair	Coordination number	Distance (nm)	fwhm (nm)	Total coordination number	Mean distance (nm)
27 mole% Ga <sub>2</sub> O <sub>3</sub> glass					
P-O	2.05*	0.1495(10)	0.012(2)	4.0*	0.1542(10)
	1.95*	0.1592(10)	0.012(2)		
Ga-O	4.05(10)	0.1890(10)	0.016(3)	6.0(2)	0.193(2)
	1.95(10)	0.2015(30)	0.022(4)		
O-O	5.2(3)	0.2520(20)	0.025(3)		
38 mole% Ga <sub>2</sub> O <sub>3</sub> glass					
P-O	2.9*	0.1515(10)	0.010(2)	4.0*	0.1540(10)
	1.1*	0.1605(10)	0.012(2)		
Ga-O	4.15(10)	0.1853(10)	0.016(3)	4.6(2)	0.187(2)
	0.45(10)	0.2020(30)	0.020(4)		
O-O	3.5(2)	0.2508(20)	0.017(3)		

\* The compositions were chosen to maintain a P-O coordination number of four.

the model  $T(r)$  functions and also the partial  $T_{ij}(r)$ 's are compared with the experimental  $T(r)$  data.

At first, glass compositions of  $x = 0.40$  and  $0.25$  have been used [15] (Chapter 2.1). Peak fitting results in P-O coordination numbers of 4.2 and 3.9 and in Ga-O coordination numbers of 4.4 and 6.3, respectively. Since 4.2 is outside the typical uncertainty of  $N_{ij}$  of a first distance peak, when  $N_{PO} = 4.0$  is expected we renormalized the glass compositions so that  $N_{PO}$  is four in both cases and started again with Gaussian fitting. The new compositions are  $x = 0.38$  and  $0.27$ , where final total Ga-O coordination numbers of 4.6 and 6.0 are obtained. All data given above are already calculated on the basis of these new compositions which includes the functions in Figs. 1, 2 and the parameters in Table 1.

#### 4. Discussion

The quality of the diffraction results presented is excellent. This is possible due to ease of corrections and to data of high  $Q$  which are available at an intense beam of high-energy photons of a synchrotron source. Reliable information of the P-O, Ga-O and O-O first-neighbor peaks of Ga phosphate glasses is obtained without using additional constraints and assumptions. The peaks are well resolved. The differences of the P-O coordination numbers from four were attributed to uncertainties of the composition. Renormalization where  $N_{PO} = 4$  is realized with improved compo-

sitions, makes no changes of distances and minor changes of other coordination numbers. So the decrease of  $N_{GaO}$  by more than unity when the Ga<sub>2</sub>O<sub>3</sub> fraction is increased from 0.27 to 0.38 is beyond any doubt. The decrease of  $N_{GaO}$  is accompanied by an evident shortening of Ga-O distances (cf. Fig. 2). The glass with 0.27 mole% Ga<sub>2</sub>O<sub>3</sub> reveals a mean  $r_{GaO}$  of 0.193 nm, while the glass with 0.38 mole% Ga<sub>2</sub>O<sub>3</sub> has a mean  $r_{GaO}$  of 0.187 nm (cf. Table 1). The corresponding Ga-O distances of the Ga(PO<sub>3</sub>)<sub>3</sub> [7] and GaPO<sub>4</sub> [2, 3] crystals with GaO<sub>6</sub> octahedra and GaO<sub>4</sub> tetrahedra are  $\sim 0.195$  nm and  $\sim 0.180$  nm, respectively. The distance  $r_{GaO}$  of the 27 mole% glass ( $N_{GaO} = 6.0$ ) is close to that of GaO<sub>6</sub> octahedra [7], while  $r_{GaO}$  of the 38 mole% glass ( $N_{GaO} = 4.6$ ) is close to the middle of Ga-O distances of the GaO<sub>6</sub> [7] and GaO<sub>4</sub> [2, 3] polyhedra.

Commonly, the O-O peak at  $\sim 0.25$  nm found for phosphate glasses is only due to edges of the PO<sub>4</sub> tetrahedra. Thus, the number  $N_{OO}$  depends on the glass composition. O<sub>T</sub> and O<sub>B</sub> sites have three and six nearest oxygen neighbors, respectively, and their ratio changes with modifier additions.  $N_{OO}$  is calculated with  $24(1-x)/(5-2x)$  [25].  $N_{OO}$  should be four at metaphosphate composition, and it decreases to  $\sim 3.4$  at pyrophosphate composition. For the two glasses studied,  $N_{OO}$  should be close to 3.93 and 3.51, while numbers of 5.2 and 3.5 are obtained (Table 1). A significant difference is found for the  $N_{OO}$  of the glass with 27 mole% Ga<sub>2</sub>O<sub>3</sub>. It remembers to the diffraction study of an Al(PO<sub>3</sub>)<sub>3</sub> glass [26], where also a tail of the O-O peak was obtained. It was argued that the short edges of AlO<sub>6</sub> octahedra contribute to this peak. Edges of GaO<sub>6</sub> octahedra with Ga-O distances of 0.193 nm would cause a contribution to O-O correlations at  $\sim 0.27$  nm, while edges of GaO<sub>4</sub> tetrahedra with Ga-O distances of 0.180 nm have lengths of  $\sim 0.29$  nm. Only the shorter lengths of the octahedral edges can contribute to the right side of an O-O peak at  $\sim 0.25$  nm. The high fraction of GaO<sub>6</sub> octahedra in the glass with 27 mole% Ga<sub>2</sub>O<sub>3</sub> causes the difference with the calculated  $N_{OO}$ , while for the glass with 38 mole% Ga<sub>2</sub>O<sub>3</sub> a similar effect is not found.

Beyond the three first-neighbor peaks a next huge peak at  $\sim 0.32$  nm is evident in the  $T(r)$  data shown in Figure 2. This peak originates from superposition of several partial correlations. A shoulder at  $\sim 0.295$  nm in the  $T(r)$  data of the glass with 27 mole% Ga<sub>2</sub>O<sub>3</sub> is attributed to P-P correlations. This feature nearly

vanishes for the other glass, which is due to the disruption of infinitely long chains to short chains, whereby  $N_{PP}$  is reduced from two to unity. The decrease in the P-P coordination number, i. e. the change from long chains to short chain fragments,  $PO_4$  dimers, and isolated groups, is verified by IR and Raman spectroscopy [15]. Studying a Ga phosphate glass close to the pyrophosphate composition a broad distribution of chain lengths with a maximum frequency for  $PO_4$  dimers was found using high-performance liquid chromatography [26]. The main peak contribution at  $\sim 0.320$  nm comes from Ga-P correlations. Ga-Ga correlations can be neglected, which is discussed below. The calculation of the P-O-P bridging angle results in  $136^\circ$  when the lengths of  $P-O_B$  bonds together with  $r_{pp}$  of  $\sim 0.295$  nm are used. A same angle of  $137^\circ$  is obtained for P-O-Ga linkages when lengths of  $P-O_T$  and Ga-O bonds ( $r_{GaO} = 0.193$  nm) together with  $r_{GaP}$  of  $\sim 0.320$  nm are used. In case of the glass of 38 mole%  $Ga_2O_3$ , the height of the peak at 0.320 nm is somewhat reduced. The peak broadens to its left flank, which even indicates a new contribution at  $\sim 0.305$  nm. This distance is explained with the increasing number of  $GaO_4$  tetrahedra which are formed instead of  $GaO_6$  octahedra. Since  $r_{GaO}$  is shorter in tetrahedra, also shorter Ga-P distances appear.

The strong change of the Ga-O coordination number in the range from metaphosphate to pyrophosphate compositions agrees with observations which were made for other phosphate glasses [10 - 14, 16]. The change was found to be an accompanying effect of the transition from a surfeit of terminal oxygen atoms with ‘dangling’  $P-O_T$  bonds to a ‘deficit’, where Me’s have to share the  $O_T$ ’s for their coordination [10 - 14]. Thus, the number  $M_{TO}$  of  $O_T$ ’s which is available for coordination of each Me is the critical number. For three-valent Me this ratio  $M_{TO} = n(O_T)/n(Me)$  is  $(1 + 2x)/x$ . In Fig. 3 this ratio  $M_{TO}$  is compared with the Ga-O coordination numbers observed. Additionally, the  $N_{MeO}$ ’s of rare earth cations,  $RE^{3+}$ , [13, 14] and of  $Al^{3+}$  [16] are given. For the  $RE^{3+}$  the minimum coordination number is about six and, thus, for  $x > 0.25$  the  $N_{REO}$  don’t follow the  $M_{TO}$  line and  $REO_n$  polyhedra start to cluster. A different behavior appears for the smaller  $Me^{3+}$  such as Al or Ga, which are well accommodated in tetrahedral environments. The Me-O coordination number well follows the number  $M_{TO}$  in that range where glasses of these oxides are obtained.

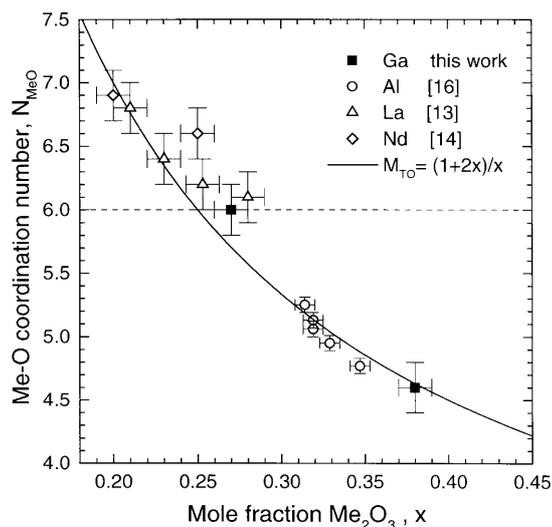


Fig. 3. Me-O coordination number vs. mole fraction of  $Me_2O_3$  in glasses  $(Me_2O_3)_x(P_2O_5)_{1-x}$  with Me = Ga [this work], La [13], Nd [14], Al [16]. The ratio  $M_{TO} = n(O_T)/n(Me)$  is given as heavy solid line. The dashed line at  $N_{MeO} = 6$  gives the lower limit of  $N_{LaO}$  and  $N_{NdO}$ , and the upper limit of  $N_{GaO}$  and  $N_{AlO}$  numbers.

The main structural features of the glasses studied where  $N_{GaO}$  is close to the  $M_{TO}$  line, i. e. close to  $n(O_T)/n(Ga)$ , are the massive formation of  $P-O_T$ -Ga bridges in addition to the  $P-O_B$ -P bridges, a lack of  $P-O_T$  bonds with  $O_T$  not having a second bond partner, a lack of  $O_T$ ’s which are shared between two  $GaO_n$  polyhedra, and also a lack of Ga-O-Ga bridges.  $N_{GaO} < M_{TO}$  would result in  $P-O_T$  bonds with  $O_T$  not having a second bond partner while  $N_{GaO} > M_{TO}$  would result in  $O_T$ ’s which are shared between  $GaO_n$  groups. Both features are unfavorable and are avoided when possible. A full transition from the above mentioned surfeit of  $O_T$ ’s to a clustering of  $MeO_n$  polyhedra due to shared  $O_T$  neighbors was found for earth-alkaline [10, 11] and RE phosphate glasses [12 - 14], but it is not detectable for Ga phosphate glasses due to the limited range of glass formation.  $P-O_T$ -Ga bridges as a main structural feature in Ga phosphate glasses have already been detected by Raman spectroscopy [15]. Also the appearance of high viscosities of the Ga phosphate melts was related [15] to the fact that oxygen atoms are mainly used in bridges, either in P-O-P or P-O-Ga links, thus all centers of structural  $PO_4$  and  $GaO_n$  units are linked with bridges.

In order to visualize which features characterize the Ga phosphate networks the structures of the Ga meta-

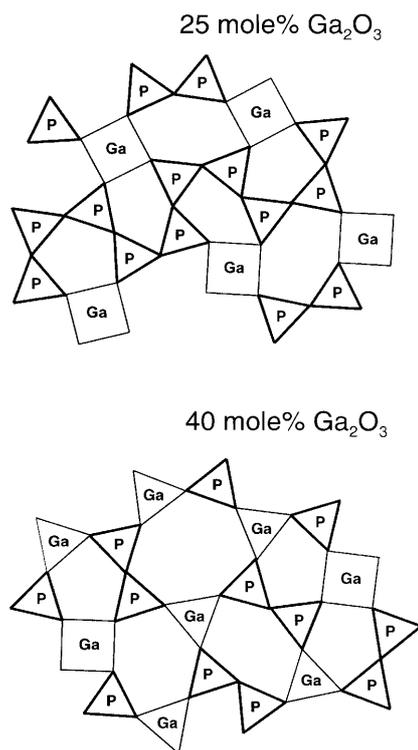


Fig. 4. Illustration of network changes from Ga meta- to pyrophosphate glasses. Triangles illustrate the tetrahedral units and squares stand for GaO<sub>6</sub> octahedra. All PO<sub>4</sub> tetrahedra link a further Ga site, all GaO<sub>4</sub> tetrahedra link a further P site and all GaO<sub>6</sub> octahedra link two further P sites vertical to the drawing plane.

and pyrophosphate glasses are illustrated in Figure 4. Due to problems in presenting three-dimensional networks in a plane, all tetrahedra are given as triangles and octahedra are given as squares. The incorporation of GaO<sub>6</sub> octahedra between highly meandering chains (25 mole% Ga<sub>2</sub>O<sub>3</sub>) results in dense arrangements of the polyhedral groups which is reflected by a high number density of atoms (80 nm<sup>-3</sup>). At pyrophosphate composition (40 mole% Ga<sub>2</sub>O<sub>3</sub>), already 75% of the Ga atoms occupy tetrahedral sites and the network expands evidently, which is expressed in a smaller number density of atoms (70 nm<sup>-3</sup>). Note, the number densities of atoms for networks formed of only tetrahedral groups whose four corners are linked with neighboring groups via bridges are even smaller. Values of ~66 nm<sup>-3</sup> are known in case of vitreous silica or of ~63 nm<sup>-3</sup> in case of vitreous germania.

The decrease of density with the change of  $N_{\text{GaO}}$  where P-O<sub>T</sub>-Ga bridges are maintained demonstrates

that the argument that the density increases when the phosphate network is disrupted from long to short chains, to PO<sub>4</sub> dimers and isolated groups is not appropriate. The compactation which is found for other polyphosphate glasses modified with, e. g., earth-alkaline and RE cations results from the ongoing process of sharing O<sub>T</sub>'s between two or three MeO<sub>n</sub> polyhedra [10, 11, 13].

It is not possible to conclude the existence of GaO<sub>5</sub> polyhedra from diffraction data while the formation of AlO<sub>5</sub> units in Al polyphosphate glasses was shown by <sup>27</sup>Al NMR [16]. When GaO<sub>5</sub> units would exist, that would seriously interfere with any crystallization. A Ga<sub>4</sub>(P<sub>2</sub>O<sub>7</sub>)<sub>3</sub> crystal, which so far is unknown, should contain 75% of Ga atoms in tetrahedra, the other part in octahedra, all that together with only PO<sub>4</sub> dimers. However, a mixture of different GaO<sub>n</sub> units stabilizes the mixture of short chains, PO<sub>4</sub> dimers and isolated PO<sub>4</sub> groups which was found in the Ga pyrophosphate glasses [15, 26]. This behavior complicates the formation of a Ga<sub>4</sub>(P<sub>2</sub>O<sub>7</sub>)<sub>3</sub> crystal and favours the separation of GaPO<sub>4</sub> phases.

## 5. Conclusions

Parameters of the P-O, Ga-O and O-O first-neighbor peaks of gallium phosphate glasses at meta- and pyrophosphate compositions are well determined using X-ray diffraction experiments of high resolving power. A strong decrease of Ga-O coordination numbers from 6.0±0.2 to 4.6±0.2 upon Ga<sub>2</sub>O<sub>3</sub> additions is detected, which is accompanied by clear shortening of Ga-O separations. Only GaO<sub>6</sub> octahedra exist at metaphosphate composition. Close to the pyrophosphate composition, already a majority of Ga atoms (75%) is found in tetrahedral sites. The Ga-O coordination number behaves equivalently with the ratio  $M_{\text{TO}} = n(\text{O}_T)/n(\text{Ga})$ , thus, with the number of terminal oxygen atoms (O<sub>T</sub>) in P-O<sub>T</sub> bonds which are available for coordination of each Ga atom. This relation leads to conclude that P-O<sub>T</sub>-Ga bridges are formed for all O<sub>T</sub> atoms. The shorter the formed phosphate chains, the more the network expands. This strange behavior is due to the change from GaO<sub>6</sub> octahedra to GaO<sub>4</sub> tetrahedra as the bridging units. The GaO<sub>n</sub> polyhedra neither share O<sub>T</sub> atoms nor form Ga-O-Ga bridges.

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